

# Interface effects of strange quark matter

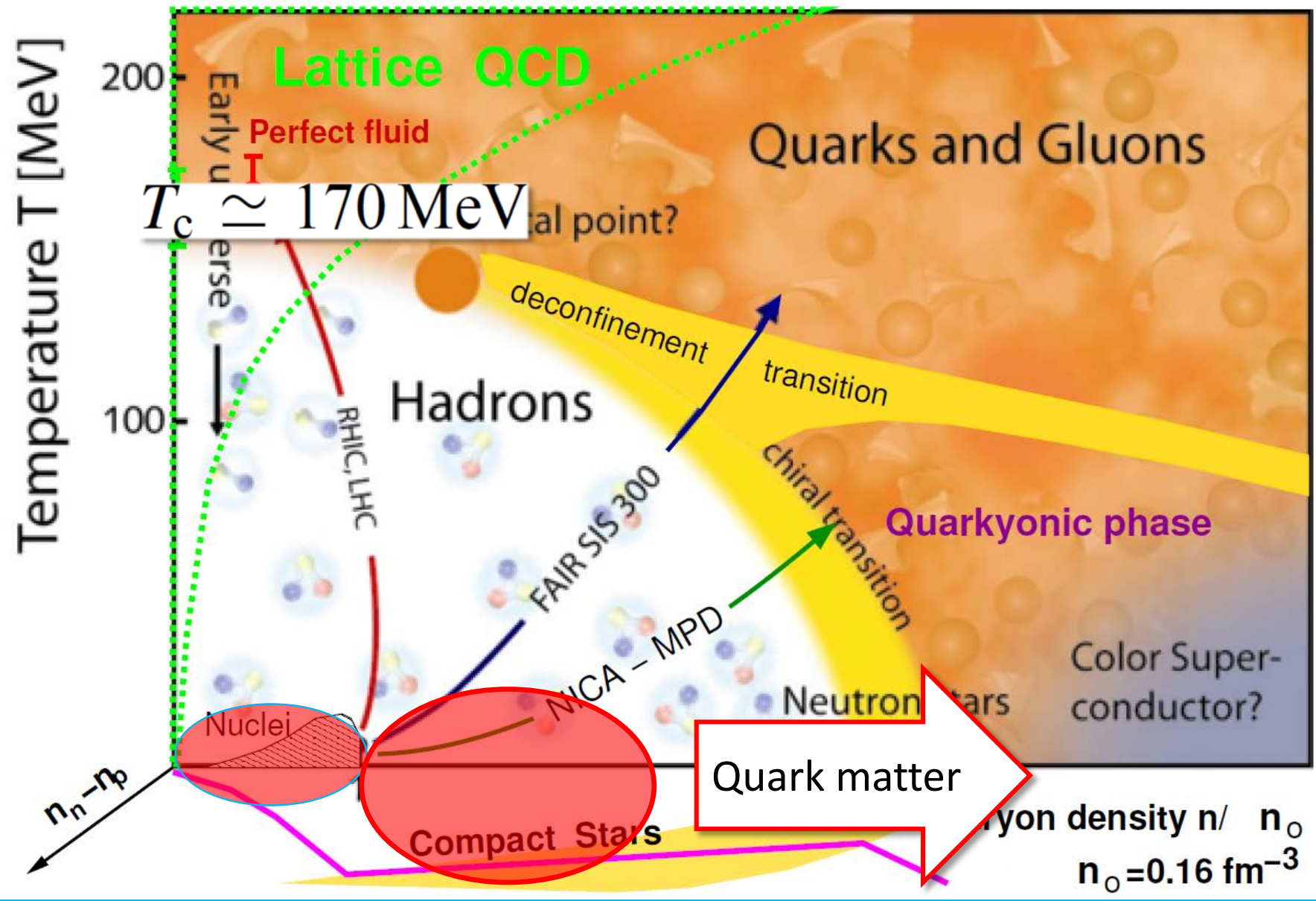
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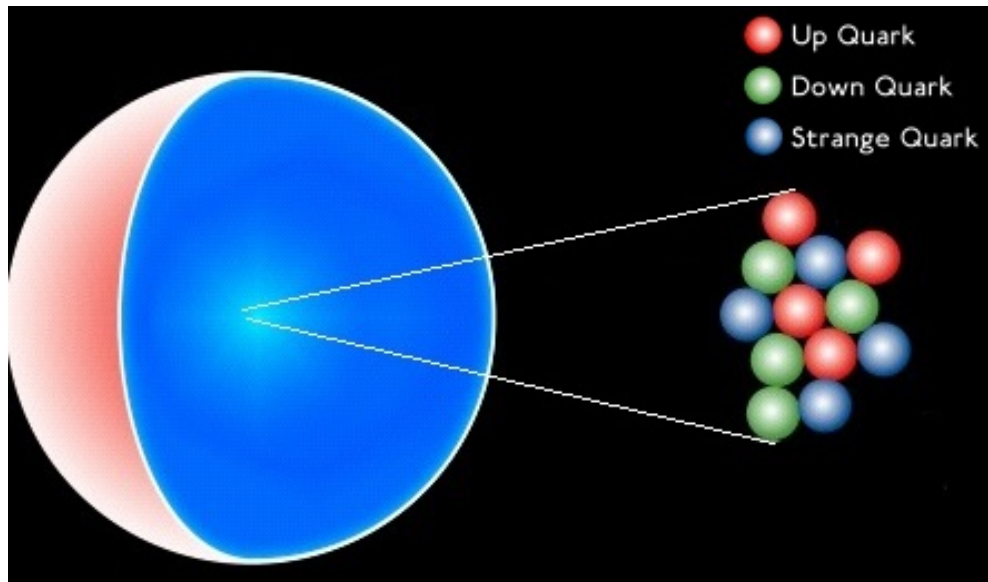
Collaborators:

Guang-Xiong Peng, En-Guang Zhao, Shan-Gui Zhou, Ting-Ting Sun,  
Wan-Lei Guo, Ding-Hui Lu, Prashanth Jaikumar

# QCD Phase Diagram [McLerran2009\_NPB195-275]



# Strange quark matter (SQM)



## Witten-Bodmer hypothesis

- Bodmer first suggested a low energy nuclear state called “**collapsed nuclei**” [Bodmer1971\_PRD4-1601];
- Witten reported on **the stability of SQM** consisting of approximately equal numbers of  $u$ ,  $d$  and  $s$  quarks, suggesting that SQM could indeed be stable even at zero external pressure [Witten1984\_PRD30-272].

If Witten-Bodmer hypothesis is true, there exists **stable lumps of SQM** with the baryon number  $A \approx 2 \sim 10^{57}$ :

### Strangelets ( $A < 10^7$ )

Comparing with **nuclei**, strangelets have: lower **charge-to-mass** ratio; larger **mass**; smaller **radius**; **spherical shape**; ...

### Nuclearites

[Rujula\_Glashow1984\_Nature312-734];

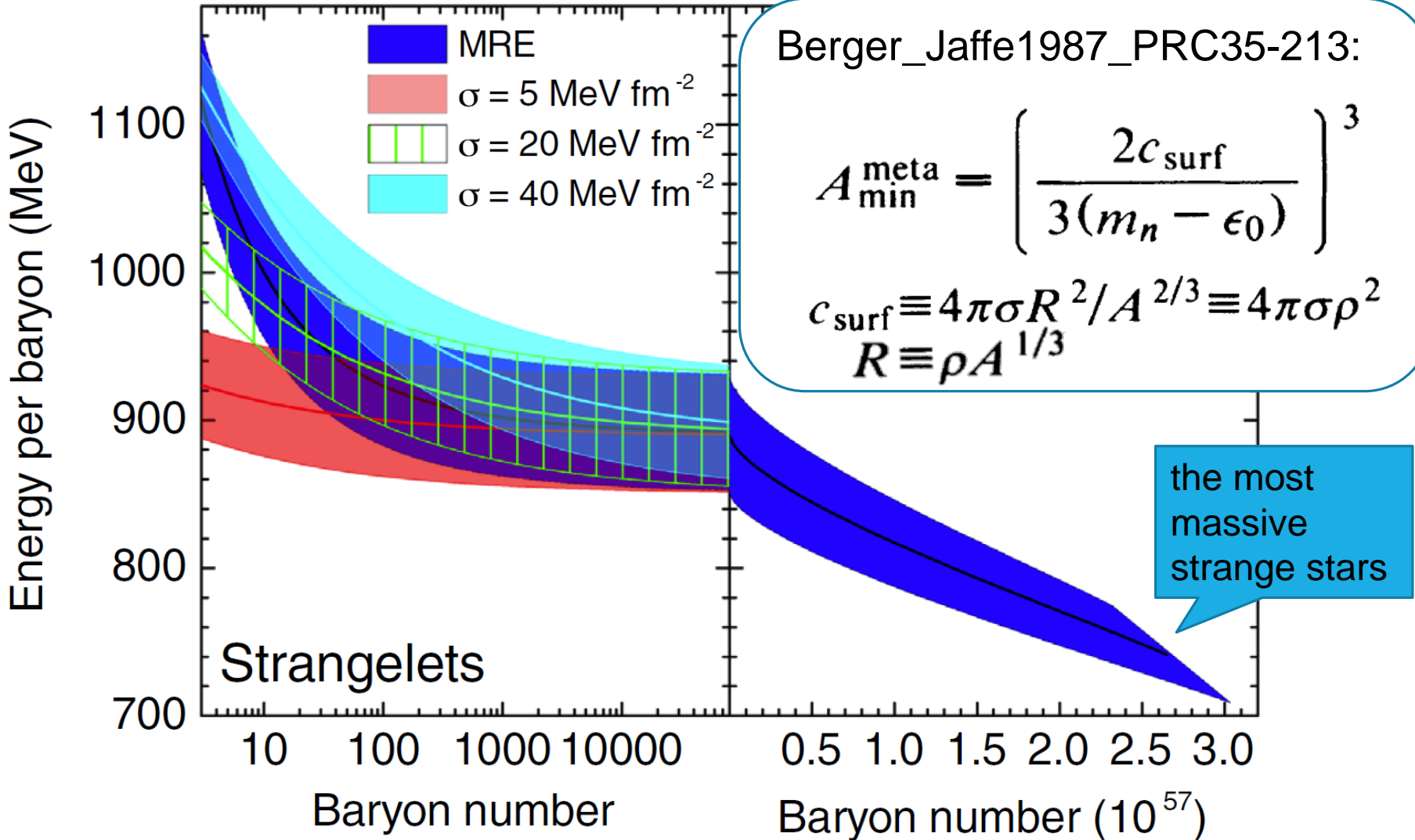
### Meteorlike Compact Ultradense Objects (CUDO)

[Rafelski\_Labun\_Birrell2013\_PRL110-111102]; ...

### Strange stars ( $A \approx 10^{57}$ )

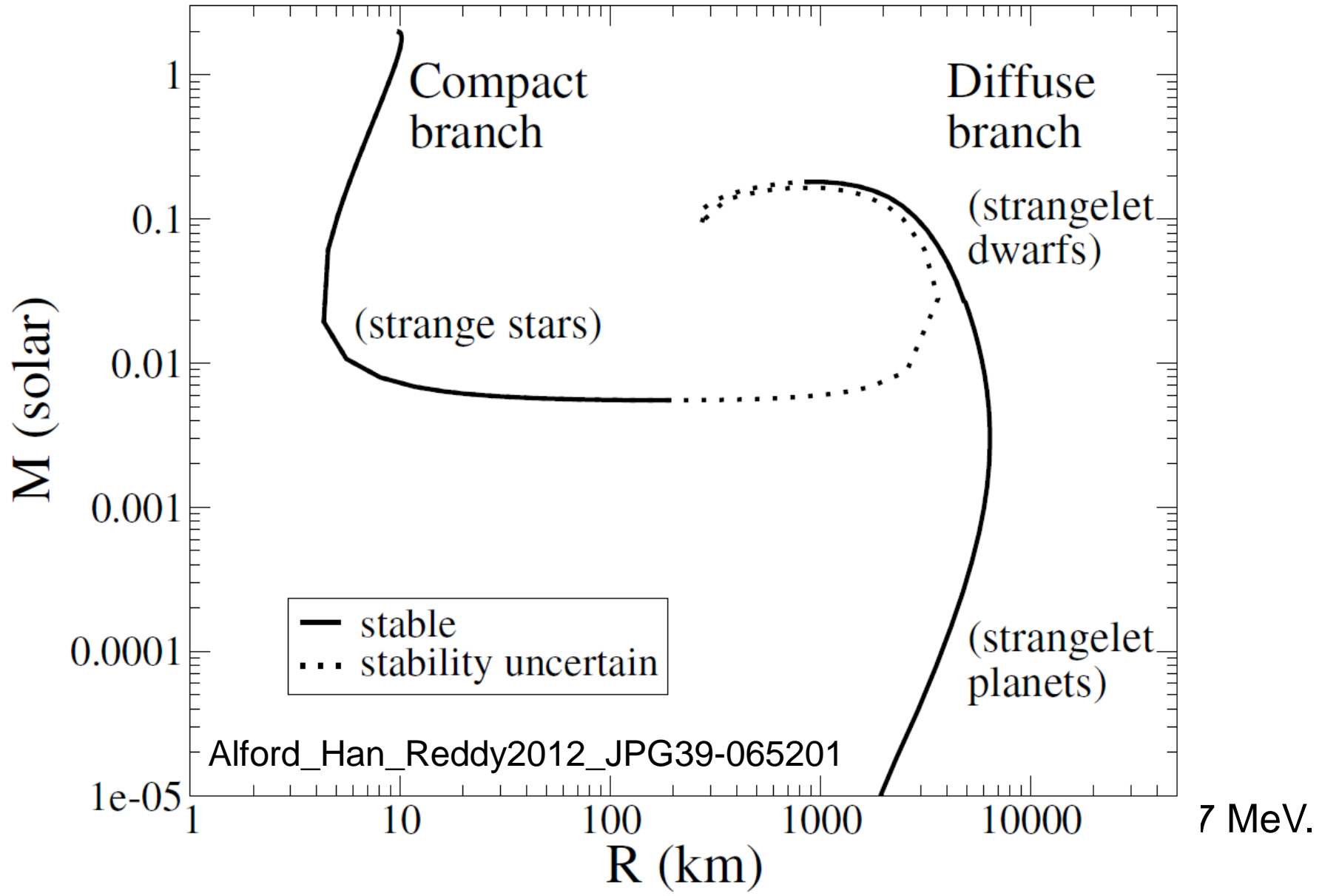
Comparing with traditional **neutron stars**, strange stars have: no **crust**; different **mass-radius** relations; smaller **radii**; higher **rotational frequencies**; ...

# Energy per baryon



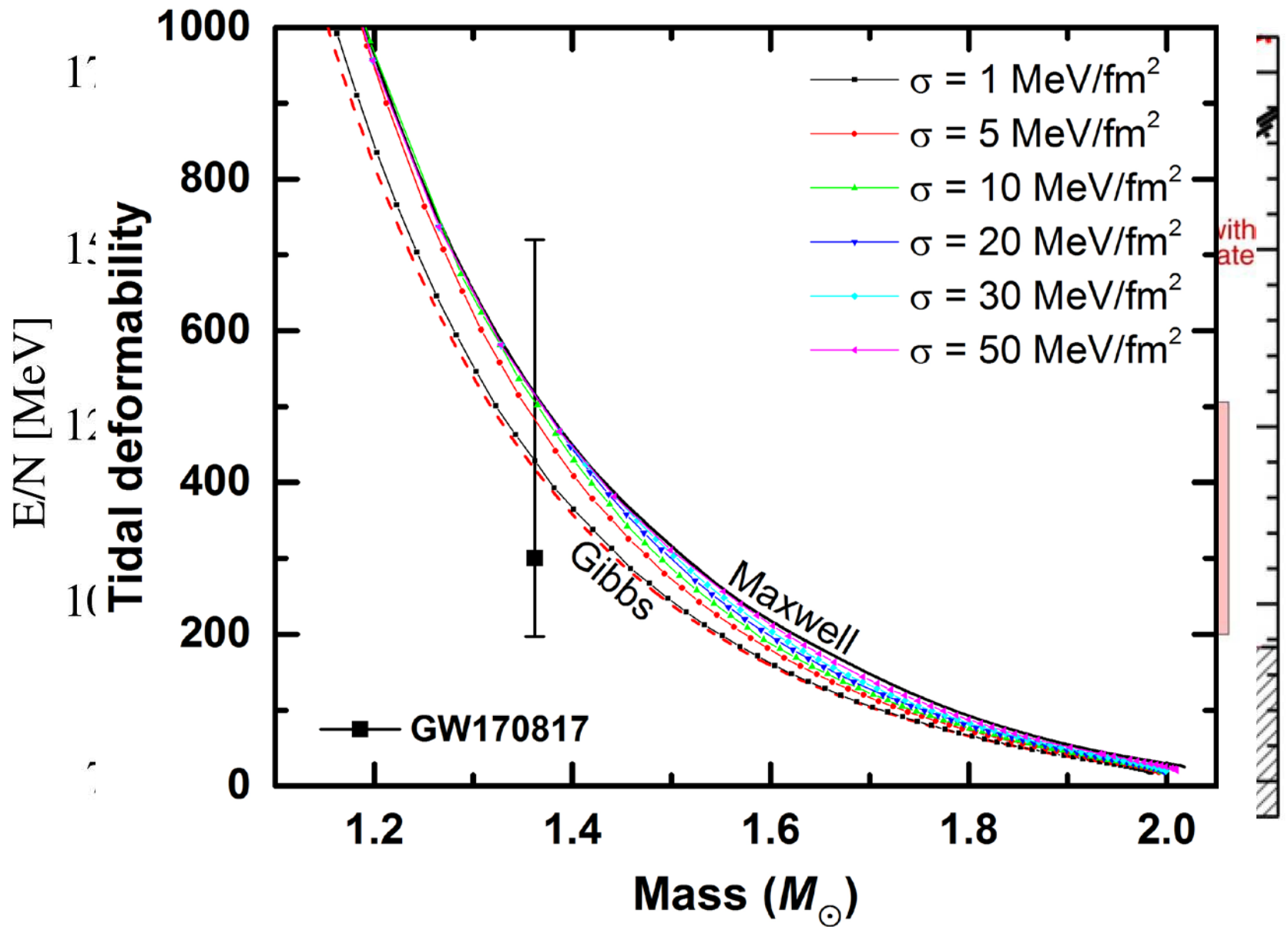
The **shaded region** correspond to the results obtained with  $B^{1/4} = 152 \pm 7 \text{ MeV}$ .  
 [Xia\_Peng\_Zhao\_Zhou2016\_PRD93-085025]

# Energy excess per baryon



Th  
[Xi]

# Unstable SQM



# Estimations of surface tension

**Lattice QCD**: Huang, Potvin, Rebbi, Sanielevici, Alves, Brower, de Forcrand, Lucini, Vettorazzo, et al.

For **vanishing** chemical potentials!

Effective models:

**Linear sigma model** [Palhares\_Fraga2010\_PRD82-125018, Pinto\_Koch\_Randrup2012\_PRC86-025203, Kroff\_Fraga2015\_PRD91-025017], **Nambu-Jona-Lasinio (NJL) model** [Garcia\_Pinto2013\_PRC88-025207, Ke\_Liu2014\_PRD89-074041], **three-flavor Polyakov-quark-meson model** [Mintz\_Stiele\_Ramos\_Schaffner-Bielich2013\_PRD87-036004], and **Dyson-Schwinger equation approach** [Gao\_Liu2016\_PRD94-094030]

$$\sigma = 5\sim 30 \text{ MeV/fm}^2$$

**Quasiparticle mode** [Wen\_Li\_Liang\_Peng2010\_PRC82-025809]

$$\sigma = 30\sim 70 \text{ MeV/fm}^2$$

**NJL model** adopting the MRE method [Lugones\_Grunfeld\_Ajmi2013\_PRC88-045803]

$$\sigma = 145\sim 165 \text{ MeV/fm}^2$$

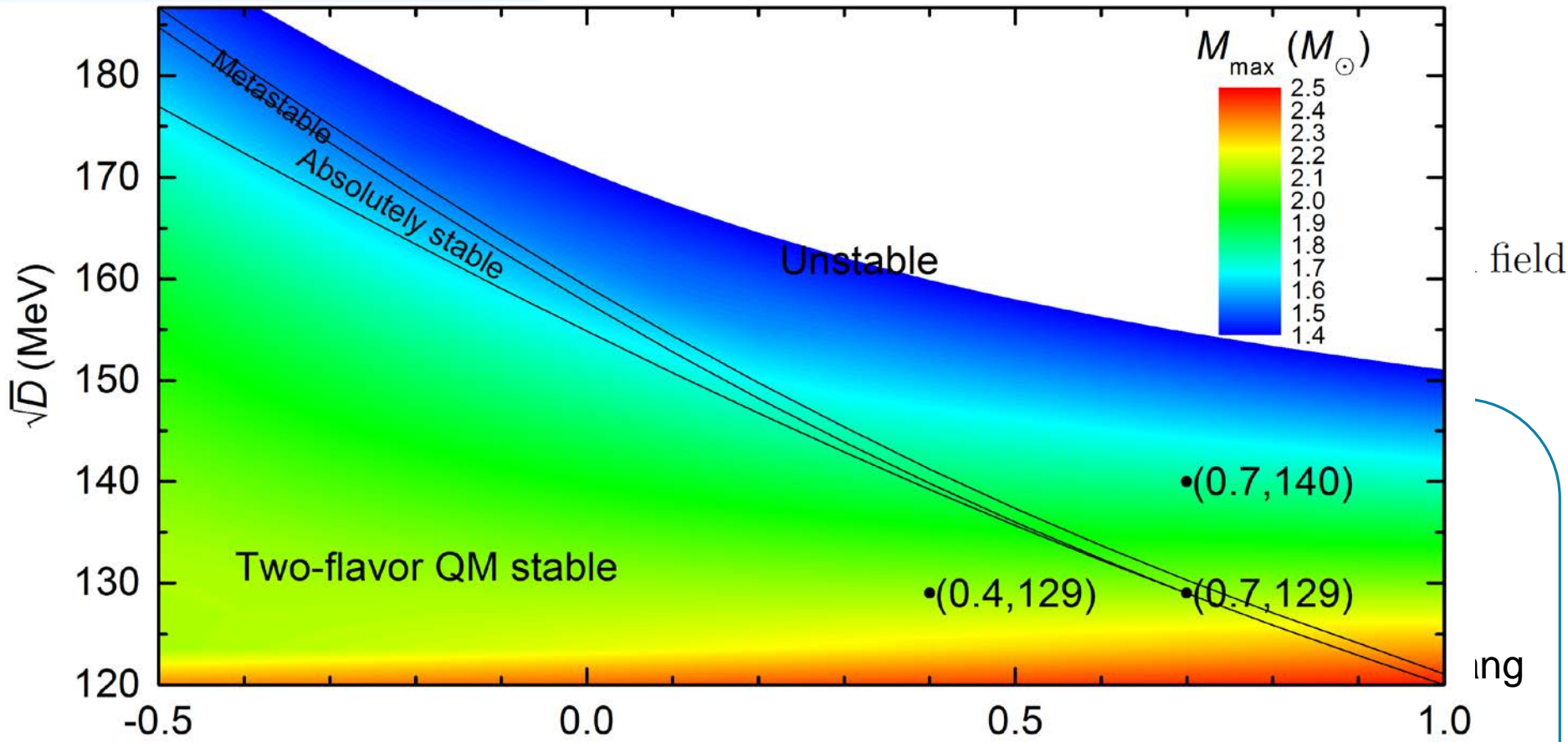
For **color-flavor locked** SQM, dimensional analysis suggests:

[Alford\_Rajagopal\_Reddy\_Wilczek2001\_PRD64-074017]

$$\sigma \approx 300 \text{ MeV/fm}^2$$

For **magnetized** SQM,  $\sigma$  has a different value in the parallel and transverse directions with respect to the magnetic field [Lugones\_Grunfeld2017\_PRC95-015804]

# Equivparticle model



- (Wen Zhong Peng, Shen Ning 2005, PRG 72-015204)
- 5 The **one-gluon-exchange interaction** was further included by Chen et al.:  
 $m_I = \frac{D}{n^{1/3}} - Cn^{1/3}$ ; (Chen, Gao, Peng 2012, CPC 36-947)
- 6 The **quark matter symmetry energy** was considered by Chu and Chen:  
 $m_I = \frac{D}{n^{1/3}} - \tau_I \delta D_I n^\alpha e^{-\beta n}$ ; (Chu, Chen 2014, ApJ 780-135)
- $C = C_1 \sqrt{\frac{\pi}{3}} \alpha$
- $C_{\max} = (3\pi^2 / N_f)^{1/3}$



# Strangelets in MFA

For spherically symmetric strangelets, the **Dirac spinor** of quarks is

$$\psi_{n\kappa m}(\mathbf{r}) = \frac{1}{r} \begin{pmatrix} iG_{n\kappa}(r) \\ F_{n\kappa}(r)\boldsymbol{\sigma} \cdot \hat{\mathbf{r}} \end{pmatrix} Y_{jm}^l(\theta, \phi)$$

radial wave functions      spinor spherical harmonics

$$\kappa = (-1)^{j+l+1/2} (j + 1/2)$$

Dirac equation

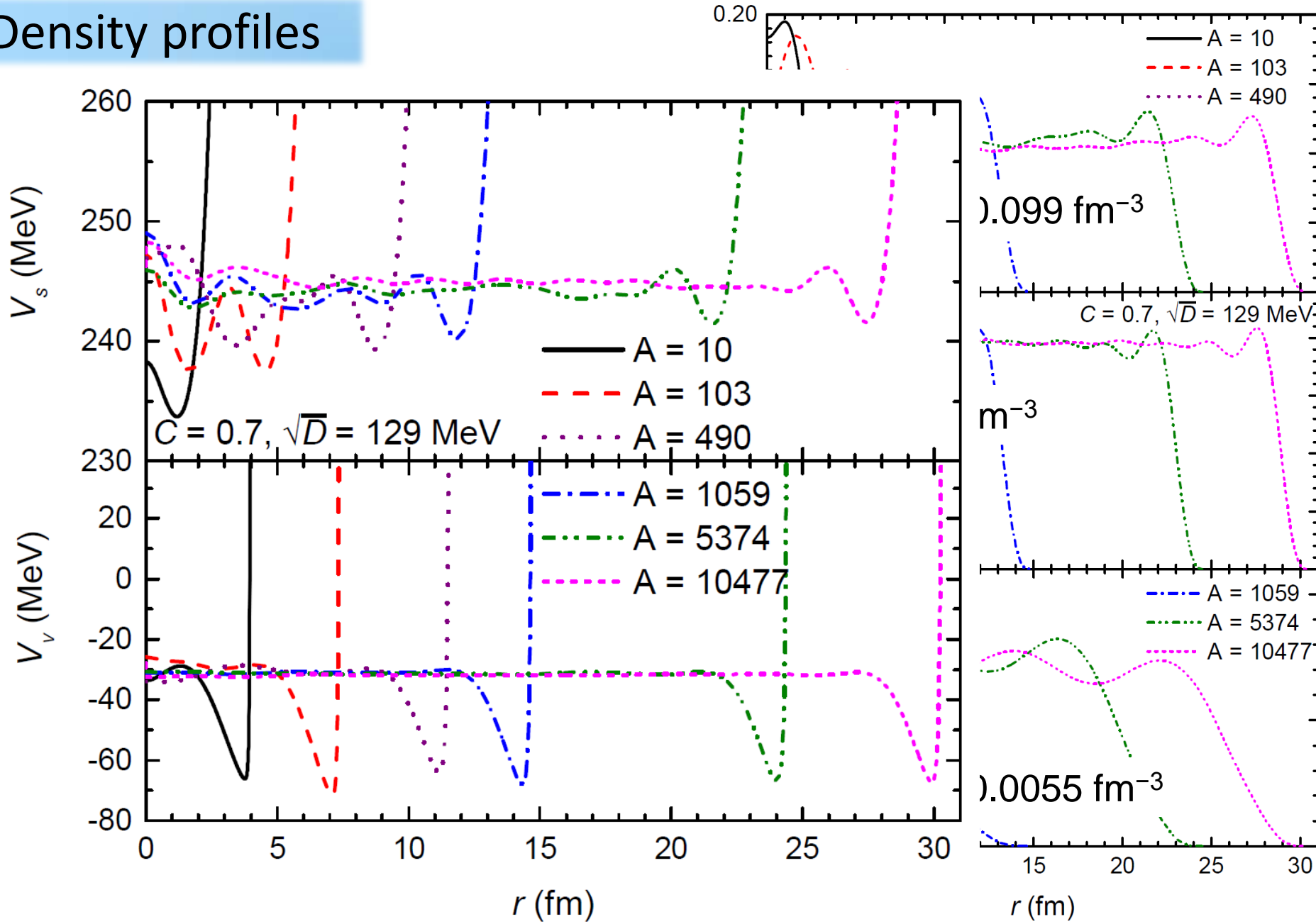
$$\begin{pmatrix} V_i + V_S & -\frac{d}{dr} + \frac{\kappa}{r} \\ \frac{d}{dr} + \frac{\kappa}{r} & V_i - V_S - 2m_{i0} \end{pmatrix} \begin{pmatrix} G_{n\kappa} \\ F_{n\kappa} \end{pmatrix} = \varepsilon_{n\kappa} \begin{pmatrix} G_{n\kappa} \\ F_{n\kappa} \end{pmatrix}$$

Mean field scalar and vector potentials

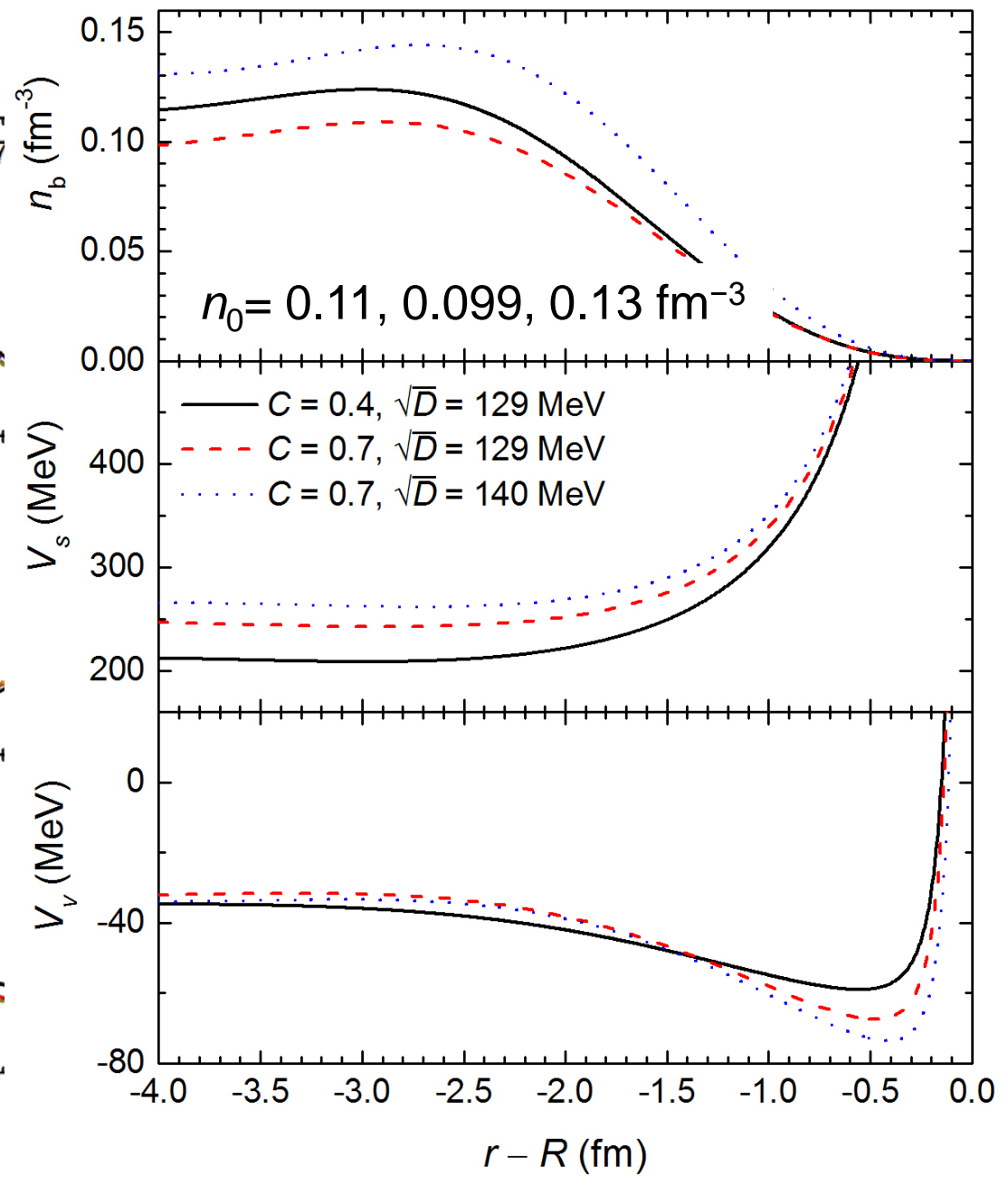
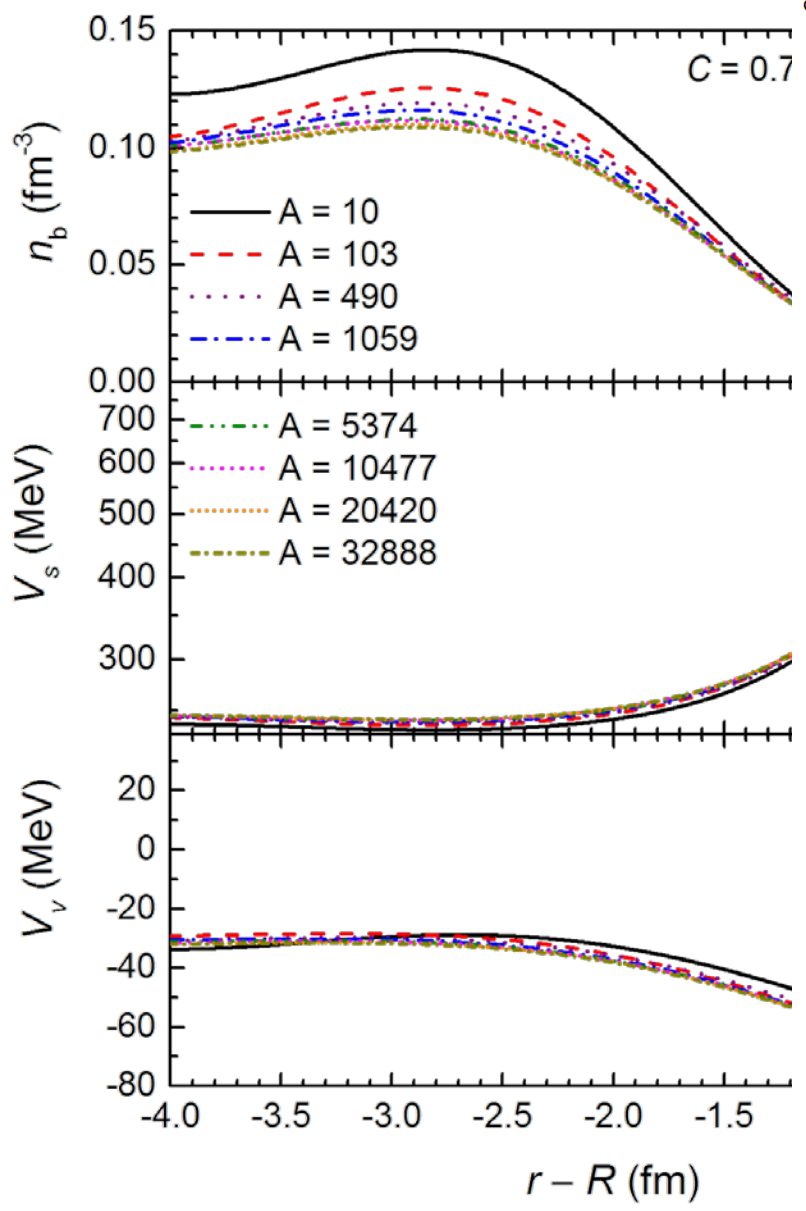
$$V_S = m_I(n_b),$$

$$V_i = \frac{1}{3} \frac{dm_I}{dn_b} \sum_{i=u,d,s} n_i^S + eq_i A_0.$$

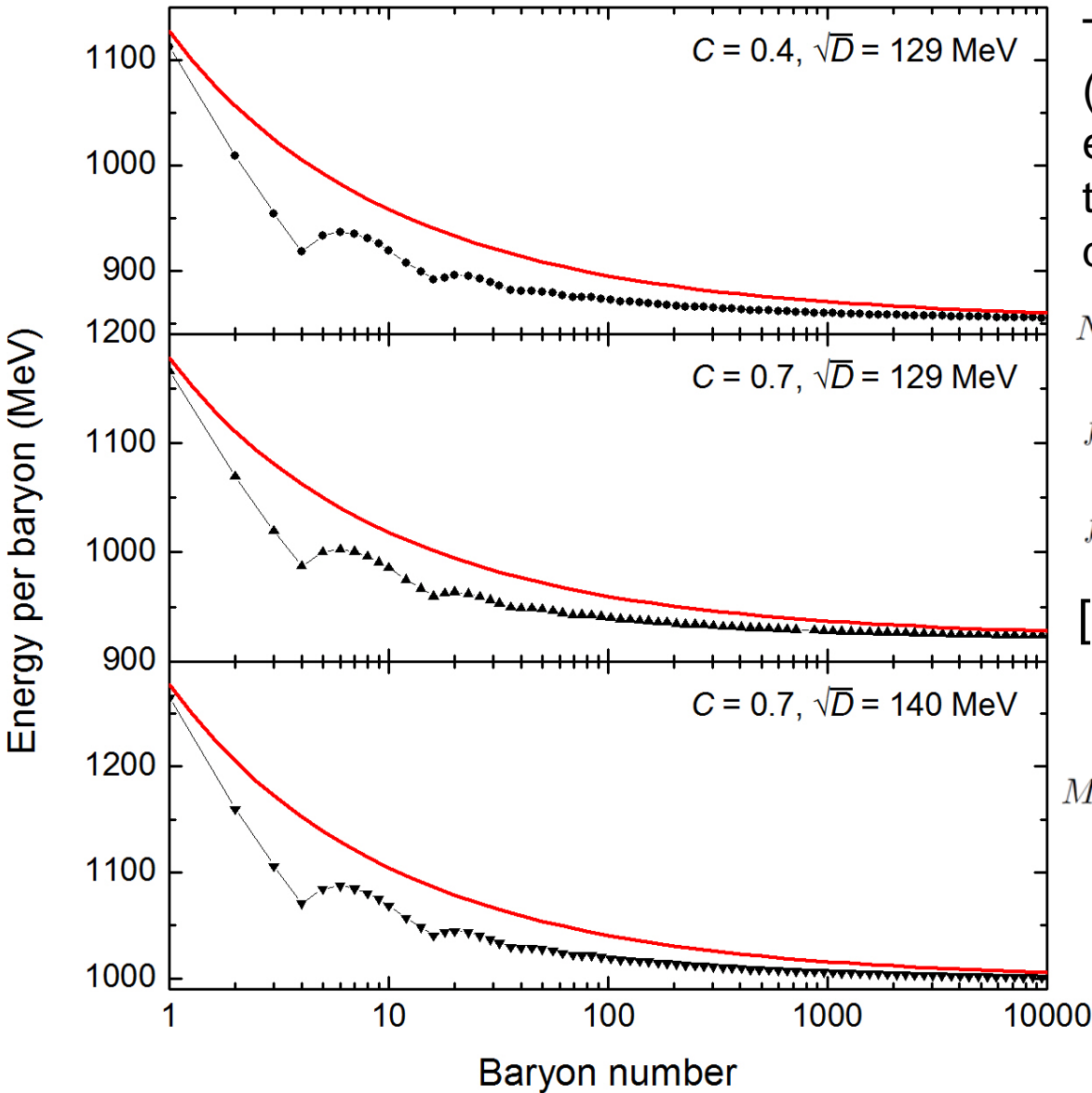
# Density profiles



# Surface structures



# Energy per baryon



The multiple reflection expansion (MRE) method: The average effects due to quark depletion are treated with a modification to the density of states, i.e.,

$$N'_i(p) = 6 \left[ \frac{p^2 v}{2\pi^2} + f_s \left( \frac{p}{m_i} \right) p s + f_c \left( \frac{p}{m_i} \right) c \right],$$

$$f_s(x) = -\frac{\eta_s}{4\pi^2} \arctan \left( \frac{1}{x} \right),$$

$$f_c(x) = \frac{\eta_c}{12\pi^2} \left[ 1 - \frac{3}{2} x \arctan \left( \frac{1}{x} \right) \right].$$

[Madsen1994\_PRD50-3328, ...]

$$N_i = \int_0^{\nu_i} N'_i(p) dp$$

$$M = \sum_{i=u,d,s} \int_0^{\nu_i} \sqrt{p^2 + m_i(n_b)^2} N'_i(p) dp + M_{ch}$$

Red solid:  $\eta_s = 1, \eta_c = 1.$

# Energy per baryon

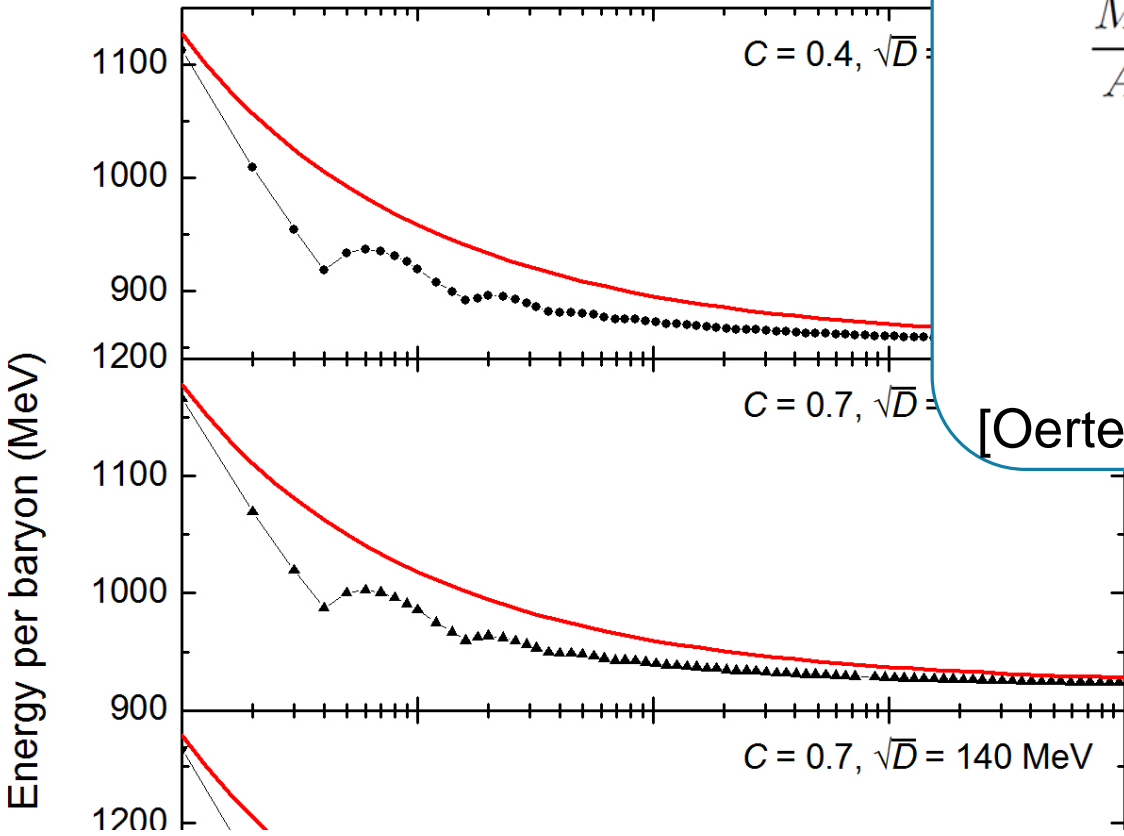
Liquid-drop type formula:

$$\frac{M}{A} = \frac{E_0}{n_0} + \frac{\alpha_S}{A^{1/3}} + \frac{\alpha_C}{A^{2/3}}$$

$$\sigma = \alpha_S \left( \frac{n_0^2}{36\pi} \right)^{1/3},$$

$$\lambda = \alpha_C \left( \frac{n_0}{384\pi^2} \right)^{1/3}.$$

[Oertel\_Urban2008\_PRD77-074015]



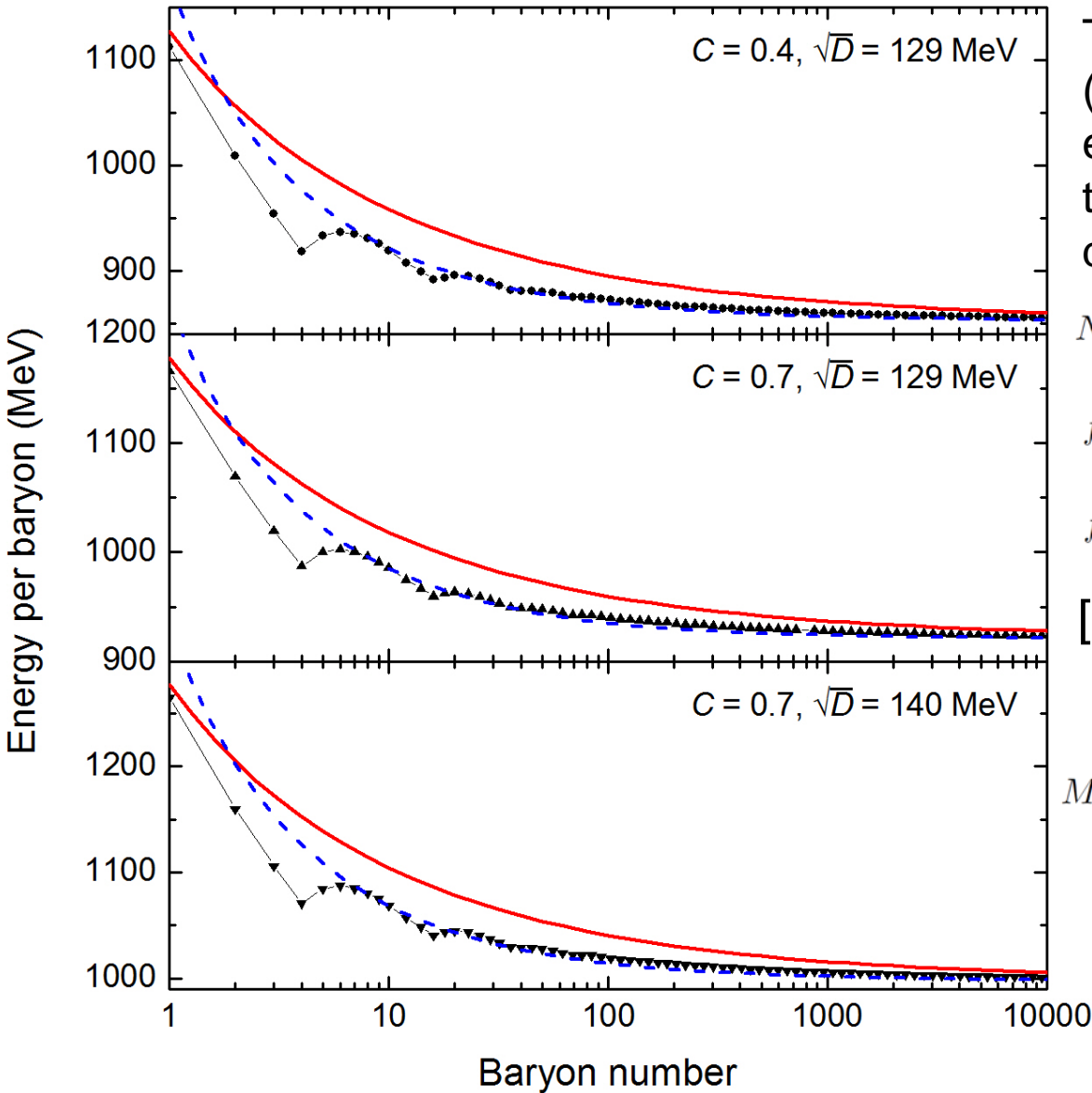
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Parameters		Bulk properties			MFA				MRE method				$\sigma^{\text{MFA}}/\sigma^{\text{MRE}}$
$C$	$\sqrt{D}$ MeV	$n_0$ $\text{fm}^{-3}$	$E_0/n_0$ MeV	$f_S$	$\alpha_S$ MeV	$\alpha_C$ MeV	$\sigma$ $\text{MeV}/\text{fm}^2$	$\lambda$ $\text{MeV}/\text{fm}$	$\alpha_S$ MeV	$\alpha_C$ MeV	$\sigma$ $\text{MeV}/\text{fm}^2$	$\lambda$ $\text{MeV}/\text{fm}$	
0.4	129	0.11	850.91	0.20	56	177	2.7	5.49	190.5	86.1	9.247	2.67	0.29
0.7	129	0.099	918.94	0.056	54	172	2.4	5.12	173.5	85.7	7.681	2.54	0.31
0.7	140	0.13	995.77	0.14	61	185	3.3	6.03	191.1	90.9	10.18	2.96	0.32

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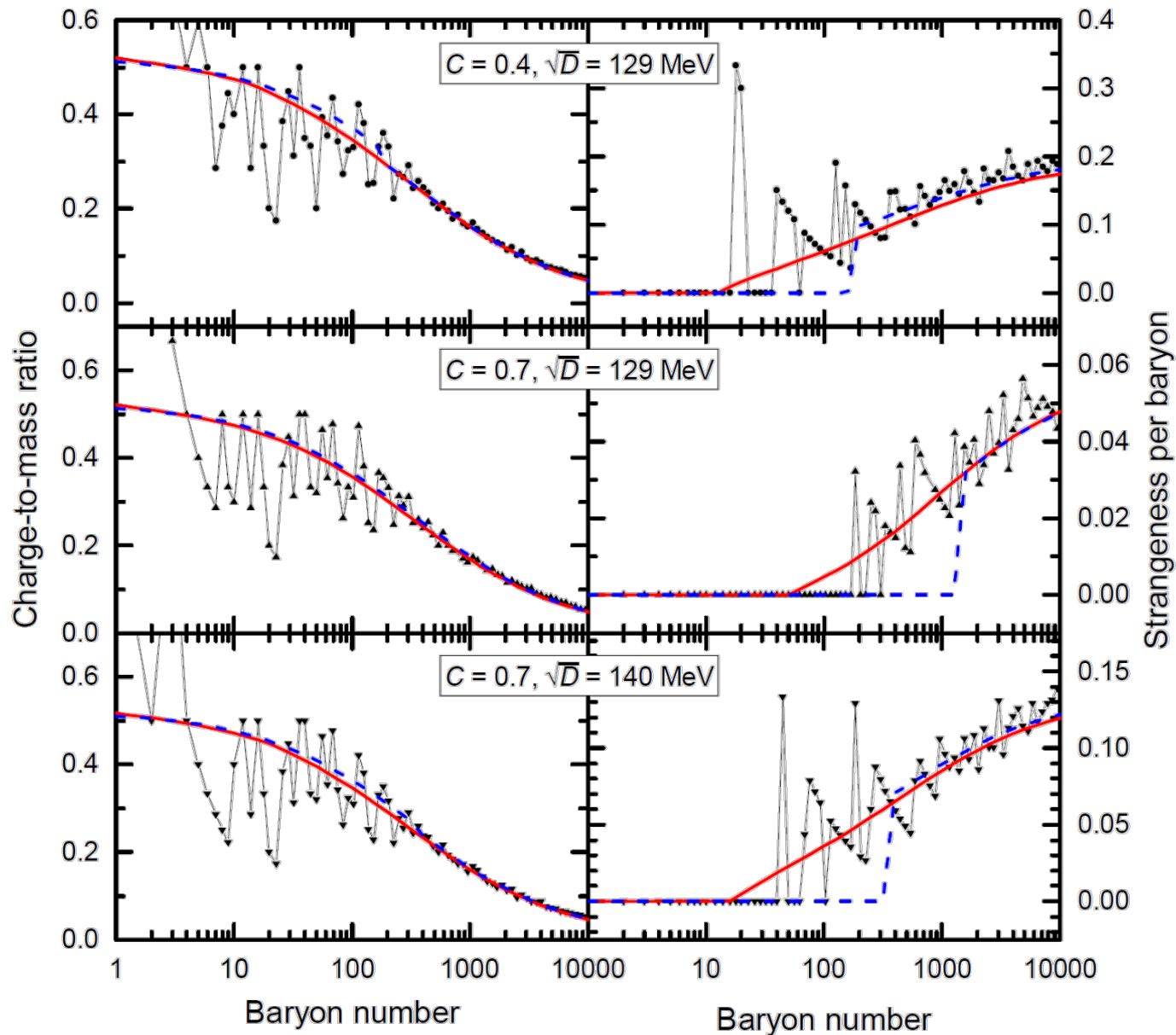
$$N_i = \int_0^{\nu_i} N'_i(p) dp$$

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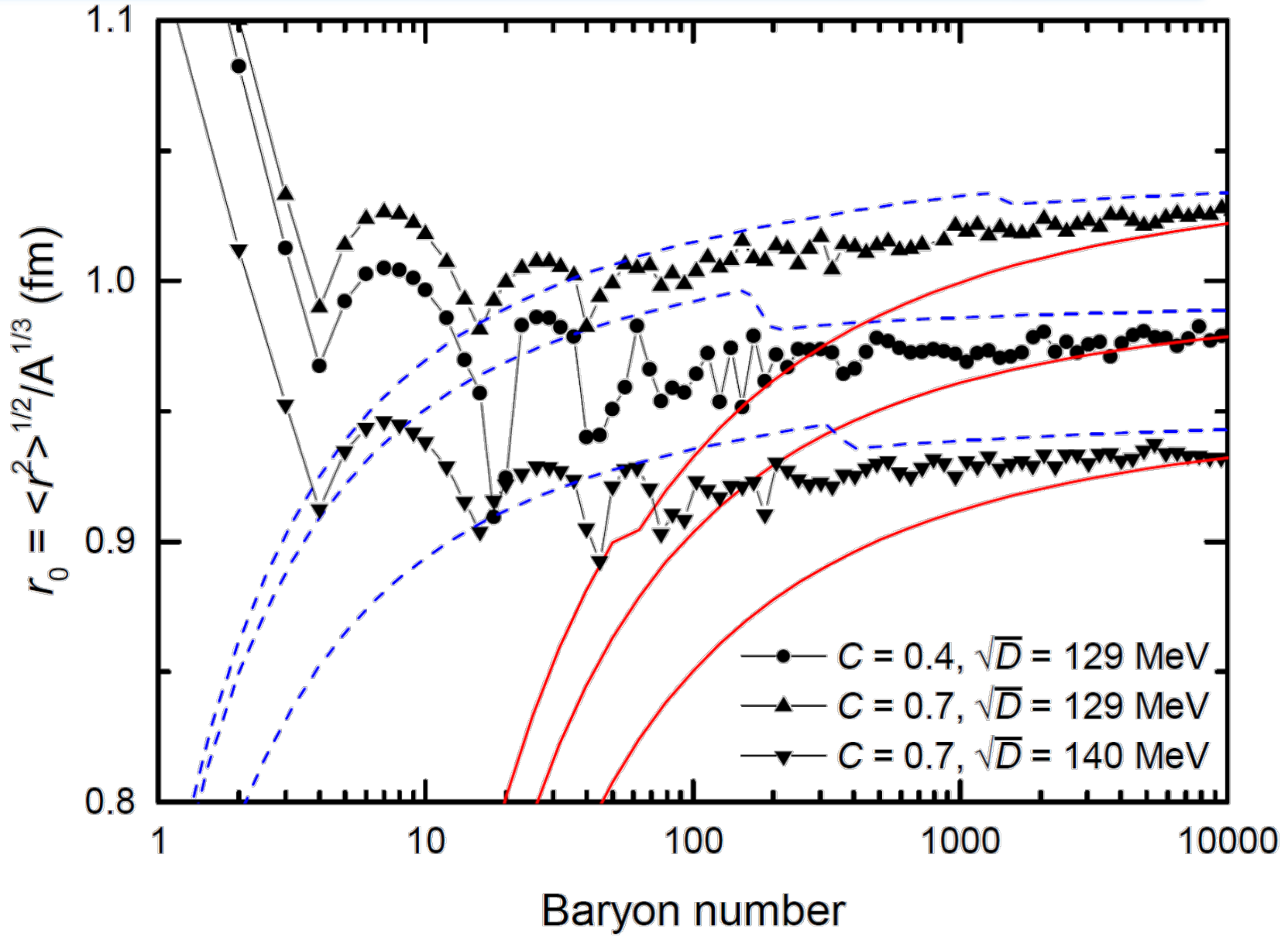
Red solid:  $\eta_s = 1, \eta_c = 1.$

Blue dashed:  $\eta_s = 0.3, \eta_c = 0.1.$

# Charge-to-mass ratio and Strangeness per baryon



# Ratio of root-mean-square radius to baryon number





# Summary

Based on an **equivparticle model**, we study the **interface effects** in strangelets adopting **mean-field approximation** (MFA). It is found that

- ① the **surface tension and curvature term** of strange quark matter (SQM) become **larger** for larger **confinement strength** and smaller **perturbative strength**;
- ② if SQM is absolutely stable and a strange star can reach 2 solar mass, the surface tension is  **$\sim 2.4 \text{ MeV}/\text{fm}^2$** ;
- ③ the **MRE** method **overestimates** the surface tension and **underestimates** the curvature term, which can be fixed by introducing proper **damping factors**; . . .

*Thank you!!!*

# The single-particle levels for u-quarks

