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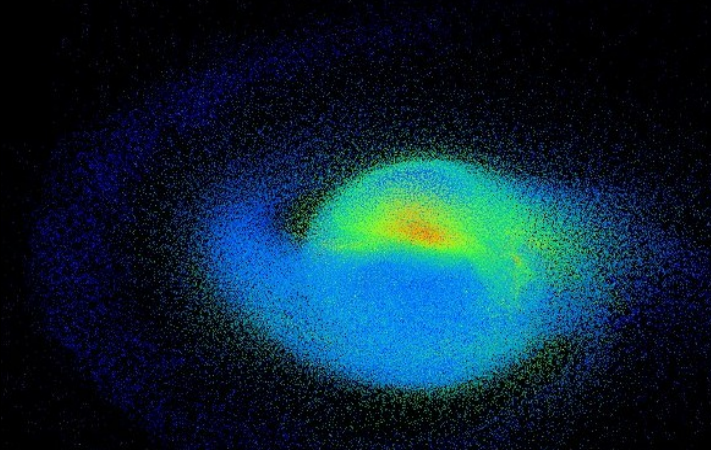
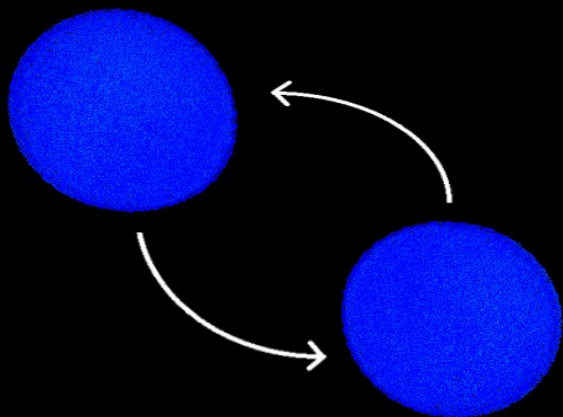
Neutron star mergers and the high-density equation of state

Xiamen-CUSTIPEN Workshop on the EOS of dense neutron-rich matter in the era of gravitational wave astronomy

Xiamen, 06/01/2019

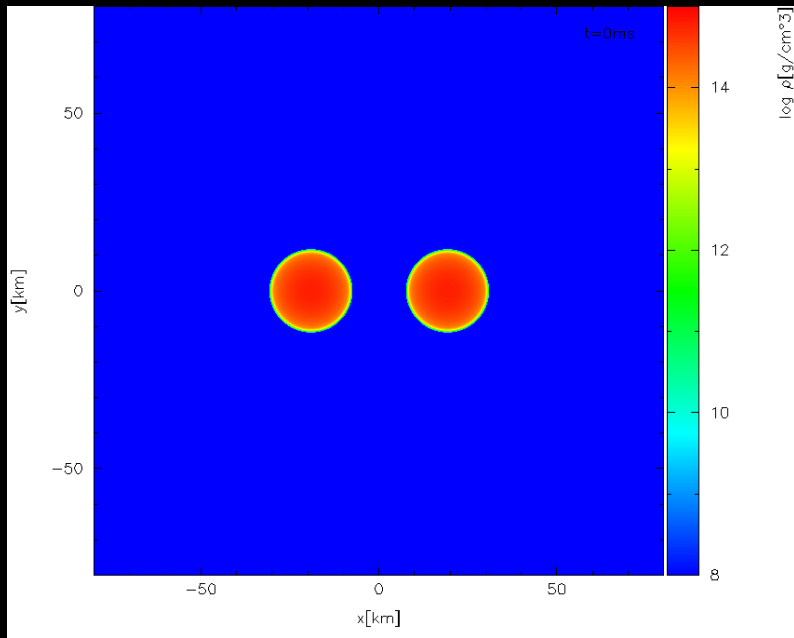
Andreas Bauswein

(GSI Darmstadt)



Outline

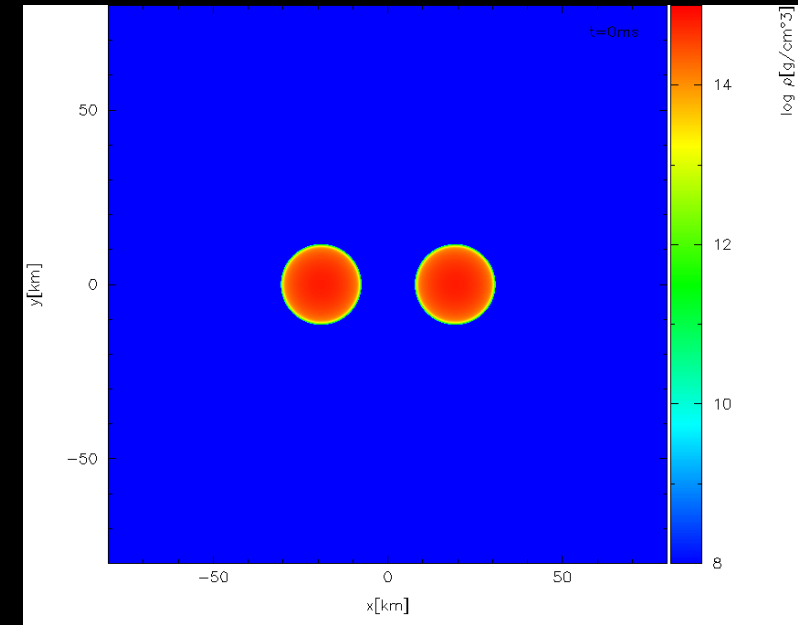
- ▶ Multi-messenger interpretation of GW170817 → lower limit on NS radii
→ Collapse behavior (EoS dependence of BH formation)
- ▶ Signatures of the QCD phase transition



$$M_{\text{tot}} = 3.4 M_{\odot}$$



$$M_{\text{tot}} = 3.5 M_{\odot}$$



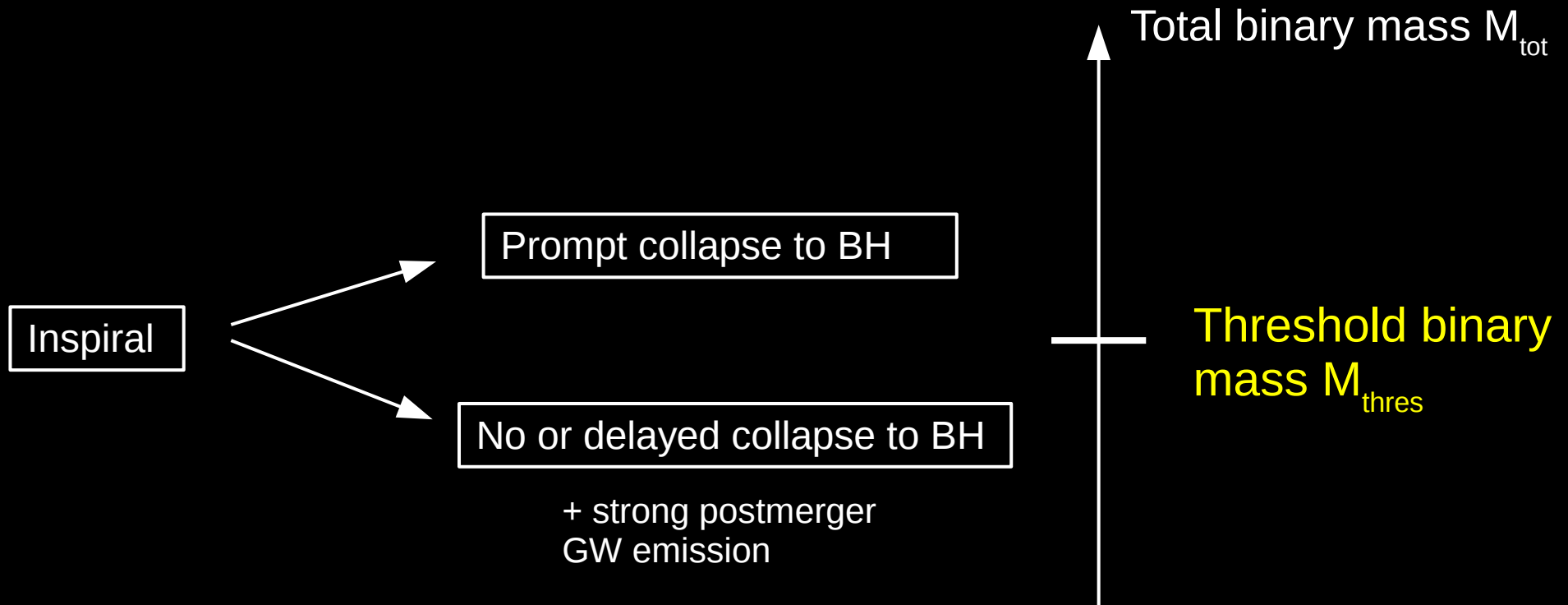
Shen EoS

$$\longrightarrow M_{\text{thres}} = (3.45 \pm 0.05) M_{\odot} \quad (\text{for this particular EoS})$$

Collapse behavior: Prompt vs. delayed (/no) BH formation

Relevant for: EoS constraints through M_{max} measurement, Conditions for short GRBs, Mass ejection, Electromagnetic counterparts powered by thermal emission, NS radius constraints !!!

Collapse behavior



EoS dependent - somehow M_{max} should play a role

Simulations reveal M_{thres}

TOV properties of nonrotating stars, i.e. EoS characteristics

Merger property from simulations

| EoS | M_{max} (M_{\odot}) | R_{max} (km) | C_{max} | $R_{1.6}$ (km) | M_{thres} (M_{\odot}) |
|-------------|-------------------------------------|--------------------------|------------------|-------------------|---------------------------------------|
| NL3 [37,38] | 2.79 | 13.43 | 0.307 | 14.81 | 3.85 |
| GS1 [39] | 2.75 | 13.27 | 0.306 | 14.79 | 3.85 |
| LS375 [40] | 2.71 | 12.34 | 0.325 | 13.71 | 3.65 |
| DD2 [38,41] | 2.42 | 11.90 | 0.300 | 13.26 | 3.35 |
| Shen [42] | 2.22 | 13.12 | 0.250 | 14.46 | 3.45 |
| TM1 [43,44] | 2.21 | 12.57 | 0.260 | 14.36 | 3.45 |
| SFHX [45] | 2.13 | 10.76 | 0.292 | 11.98 | 3.05 |
| GS2 [46] | 2.09 | 11.78 | 0.262 | 13.31 | 3.25 |
| SFHO [45] | 2.06 | 10.32 | 0.294 | 11.76 | 2.95 |
| LS220 [40] | 2.04 | 10.62 | 0.284 | 12.43 | 3.05 |
| TMA [44,47] | 2.02 | 12.09 | 0.247 | 13.73 | 3.25 |
| IUF [38,48] | 1.95 | 11.31 | 0.255 | 12.57 | 3.05 |

Bauswein et al. 2013

Smooth particle hydrodynamics + conformal flatness

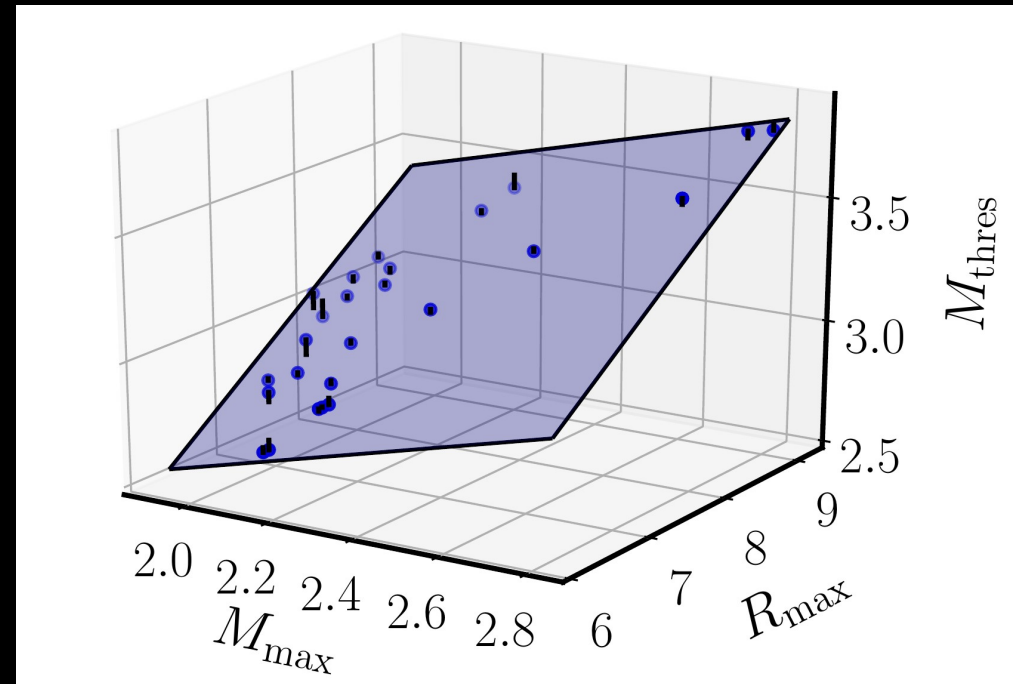
Threshold binary mass

- ▶ Empirical relation from simulations with different M_{tot} and EoS
- ▶ Fits (to good accuracy):

$$M_{\text{thres}} = M_{\text{thres}}(M_{\text{max}}, R_{\text{max}}) = \left(-3.38 \frac{GM_{\text{max}}}{c^2 R_{\text{max}}} + 2.43 \right) M_{\text{max}}$$

$$M_{\text{thres}} = M_{\text{thres}}(M_{\text{max}}, R_{1.6}) = \left(-3.6 \frac{GM_{\text{max}}}{c^2 R_{1.6}} + 2.38 \right) M_{\text{max}}$$

- ▶ Both better than $0.06 M_{\text{sun}}$



EoS constraints from GW170817*

→ lower bound on NS radii

(recall: upper bound from tidal deformability)

* See also Margalit & Metzger 2017, Shibata et al. 2017, Rezzolla et al. 2018, Radice et al. 2018, Ruiz & Shapiro 2018, ... for other EoS constraints in the context of GW170817

A simple but robust NS radius constraint from GW170817

- ▶ High ejecta mass inferred from electromagnetic transient

(high compared to simulations)

→ provides strong support for a delayed/no collapse in GW170817

→ even asymmetric mergers that directly collapse do not produce such massive ejecta

| Reference | $m_{\text{dyn}} [M_{\odot}]$ | $m_w [M_{\odot}]$ |
|-----------------------------|------------------------------|-------------------|
| Abbott et al. (2017a) | 0.001 – 0.01 | – |
| Arcavi et al. (2017) | – | 0.02 – 0.025 |
| Cowperthwaite et al. (2017) | 0.04 | 0.01 |
| Chornock et al. (2017) | 0.035 | 0.02 |
| Evans et al. (2017) | 0.002 – 0.03 | 0.03 – 0.1 |
| Kasen et al. (2017) | 0.04 | 0.025 |
| Kasliwal et al. (2017b) | > 0.02 | > 0.03 |
| Nicholl et al. (2017) | 0.03 | – |
| Perego et al. (2017) | 0.005 – 0.01 | 10^{-5} – 0.024 |
| Rosswog et al. (2017) | 0.01 | 0.03 |
| Smartt et al. (2017) | 0.03 – 0.05 | 0.018 |
| Tanaka et al. (2017a) | 0.01 | 0.03 |
| Tanvir et al. (2017) | 0.002 – 0.01 | 0.015 |
| Troja et al. (2017) | 0.001 – 0.01 | 0.015 – 0.03 |

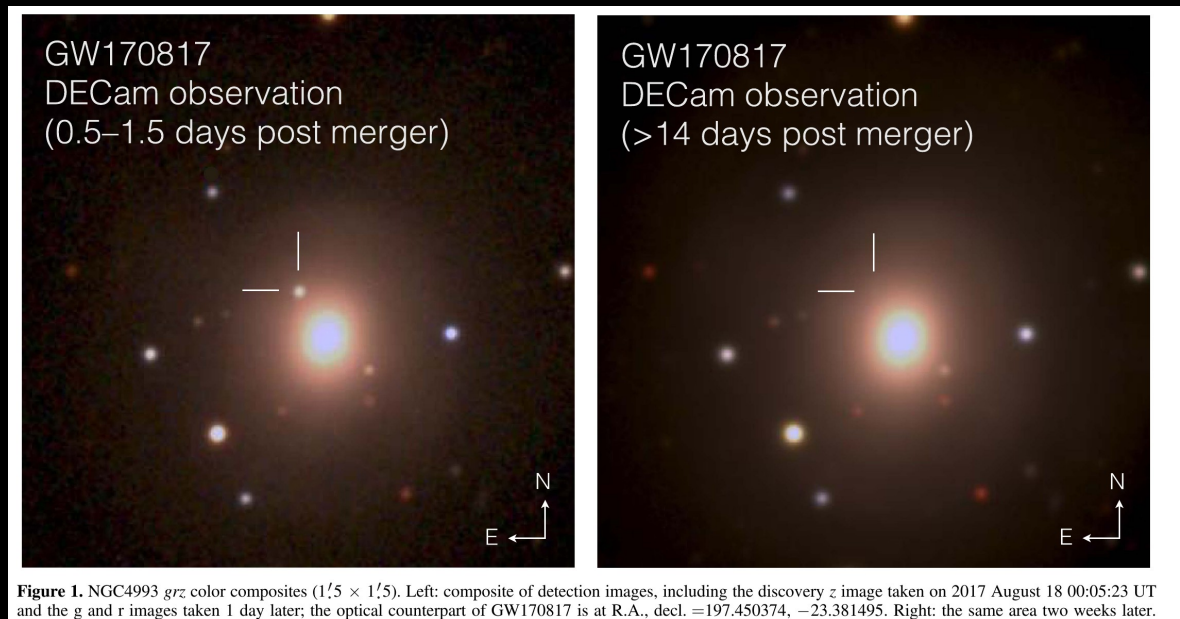
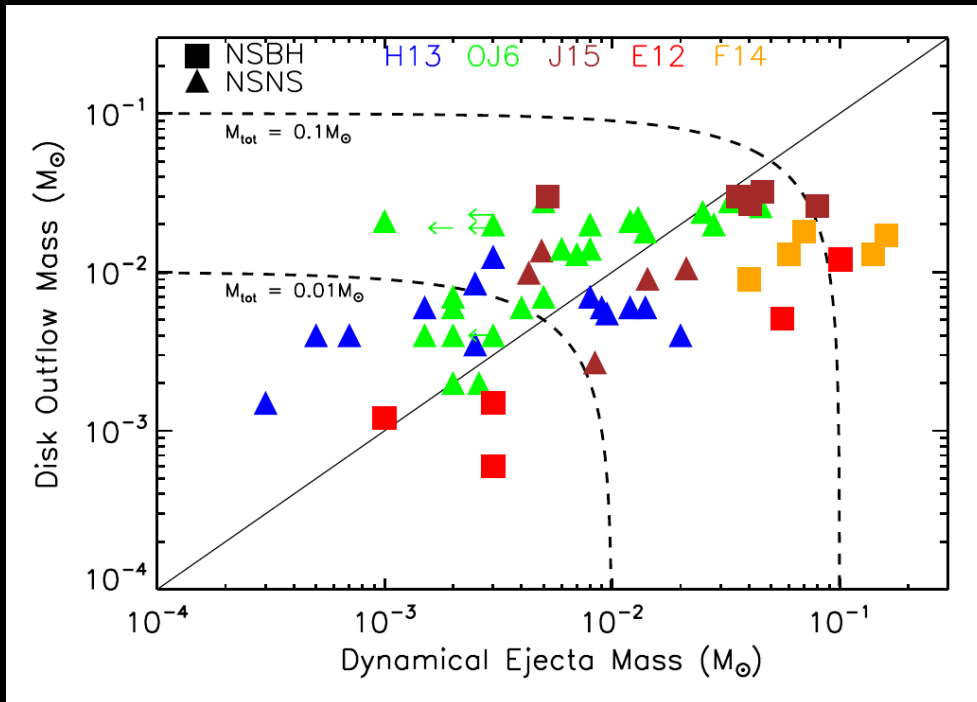


Figure 1. NGC4993 $g/r/z$ color composites ($1'5 \times 1'5$). Left: composite of detection images, including the discovery z image taken on 2017 August 18 00:05:23 UT and the g and r images taken 1 day later; the optical counterpart of GW170817 is at R.A., decl. =197.450374, -23.381495 . Right: the same area two weeks later.

Soares-Santos et al 2017

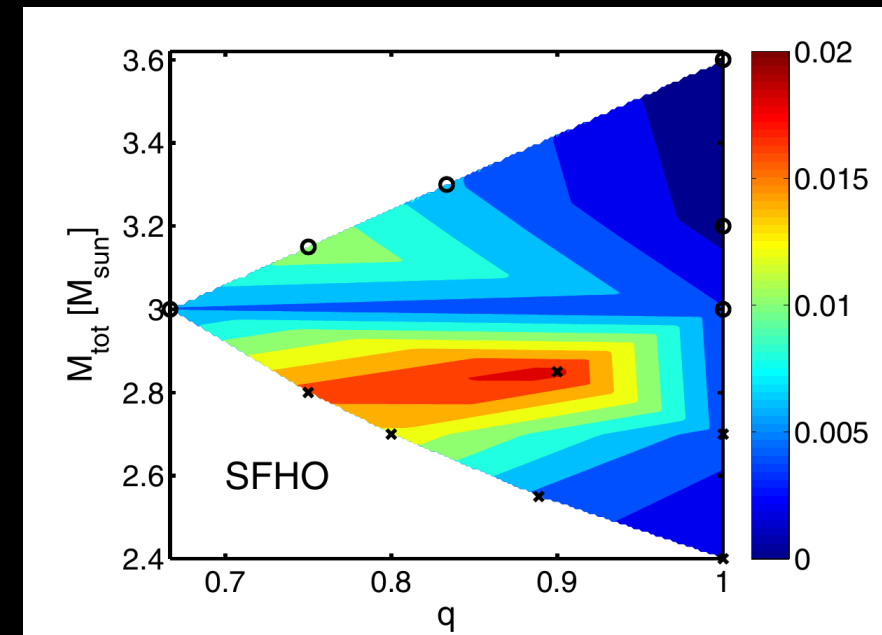
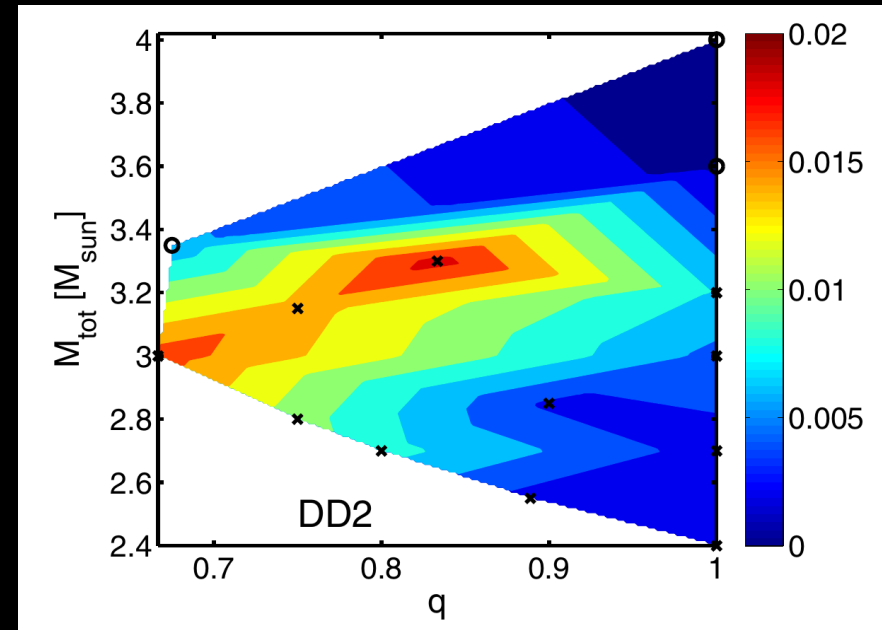
Compilation in Cote et al 2018

- ▶ Ejecta masses depend on EoS and binary masses
- ▶ Note: high mass points already to soft EoS (tentatively/qualitatively)
- ▶ Prompt collapse leads to reduced ejecta mass
- ▶ Light curve depends on ejecta mass:
→ 0.02 - 0.05 M_{sun} point to delayed collapse



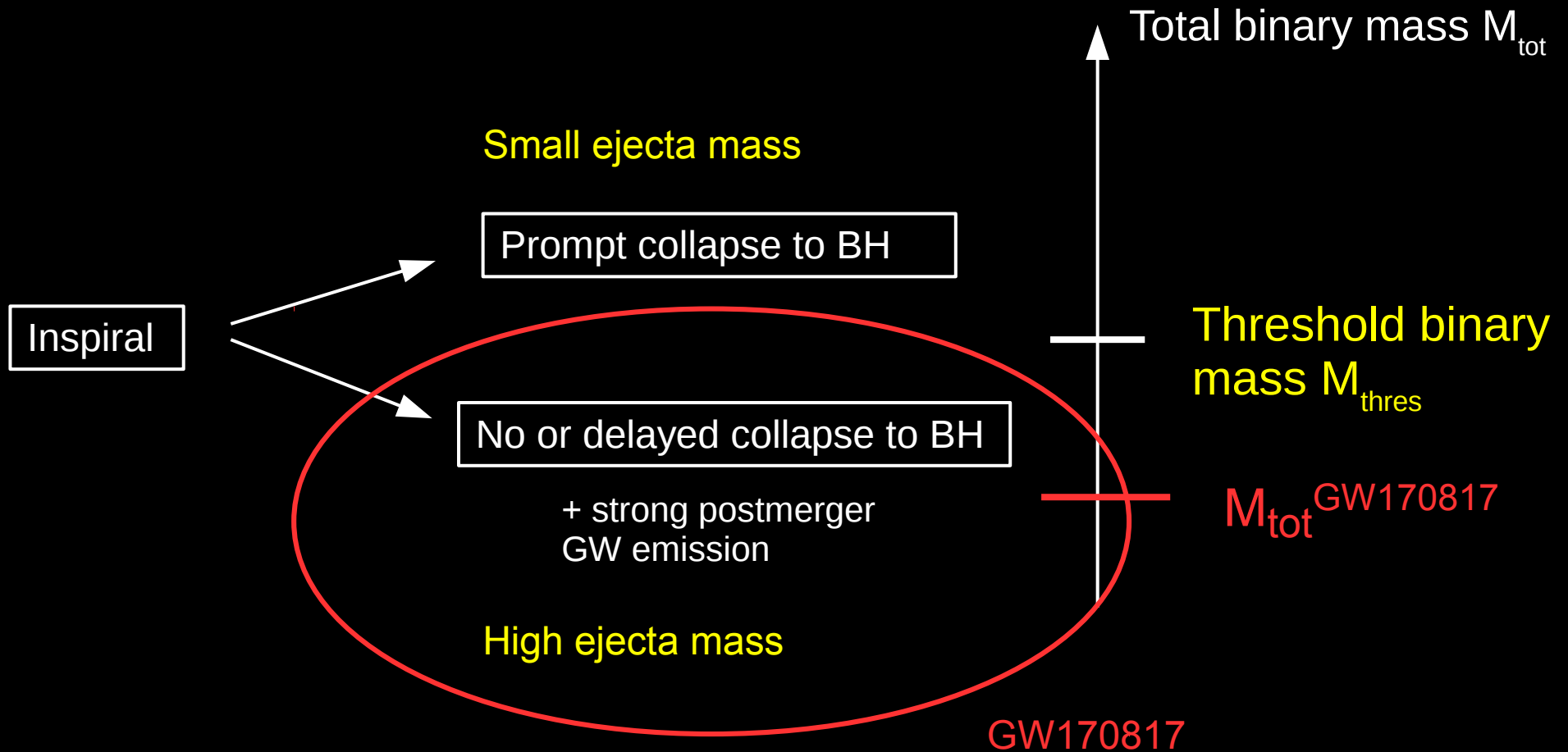
Compilation Wu et al 2016: dynamical and secular ejecta comparable

Only dynamical ejecta



Bauswein et al. 2013

Collapse behavior



(1) If GW170817 was a delayed (/no) collapse:

$$M_{\text{thres}} > M_{\text{tot}}^{\text{GW170817}}$$

(2) Recall: empirical relation for threshold binary mass for prompt collapse:

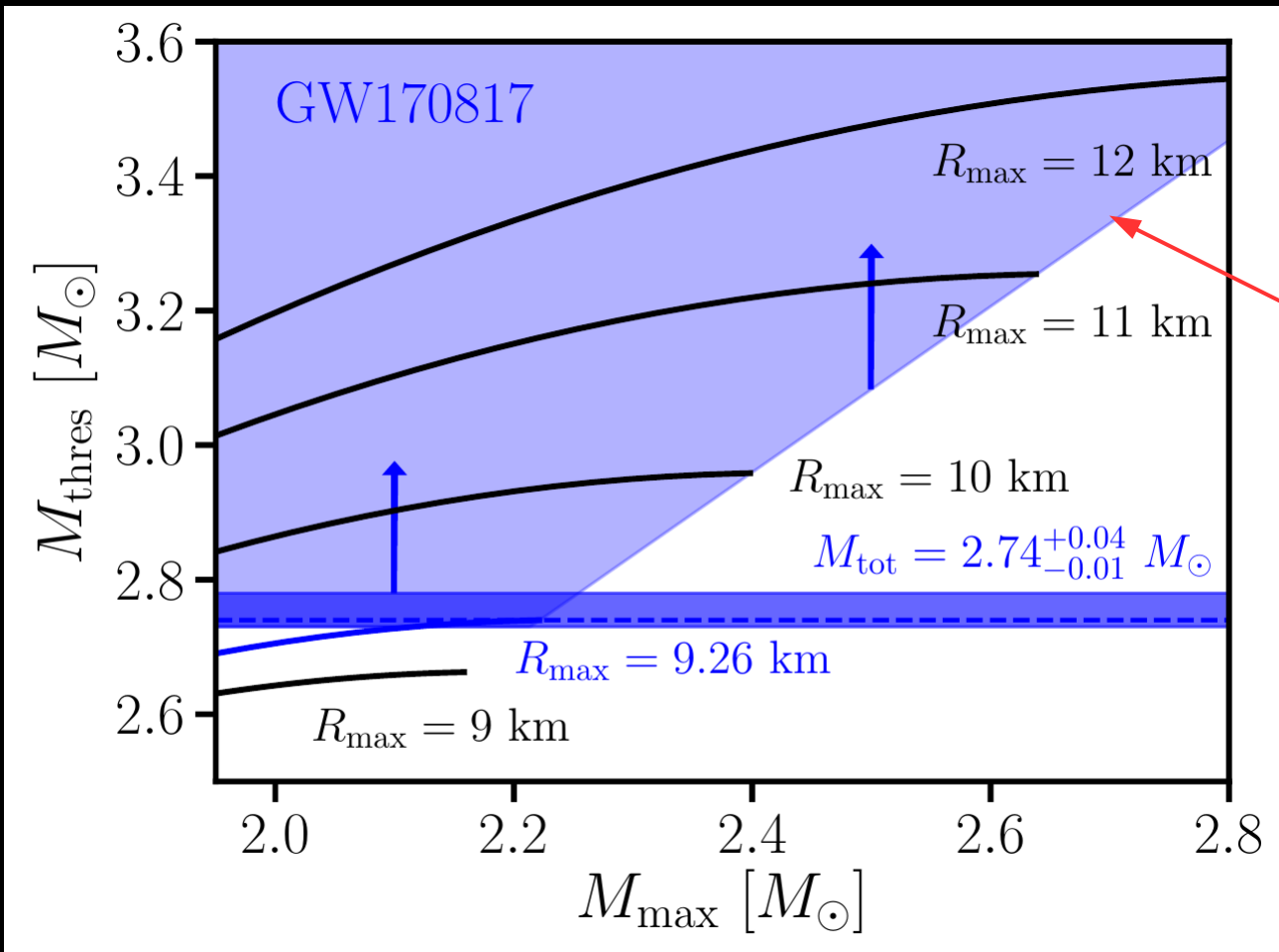
$$M_{\text{thres}} = \left(-3.38 \frac{G M_{\text{max}}}{c^2 R_{\text{max}}} + 2.43 \right) M_{\text{max}} > 2.74 M_{\odot} \quad (\text{with } M_{\text{max}}, R_{\text{max}} \text{ unknown})$$

(3) Causality: speed of sound $v_s \leq c \Rightarrow M_{\text{max}} \leq \frac{1}{2.82} \frac{c^2 R_{\text{max}}}{G}$

► Putting things together:

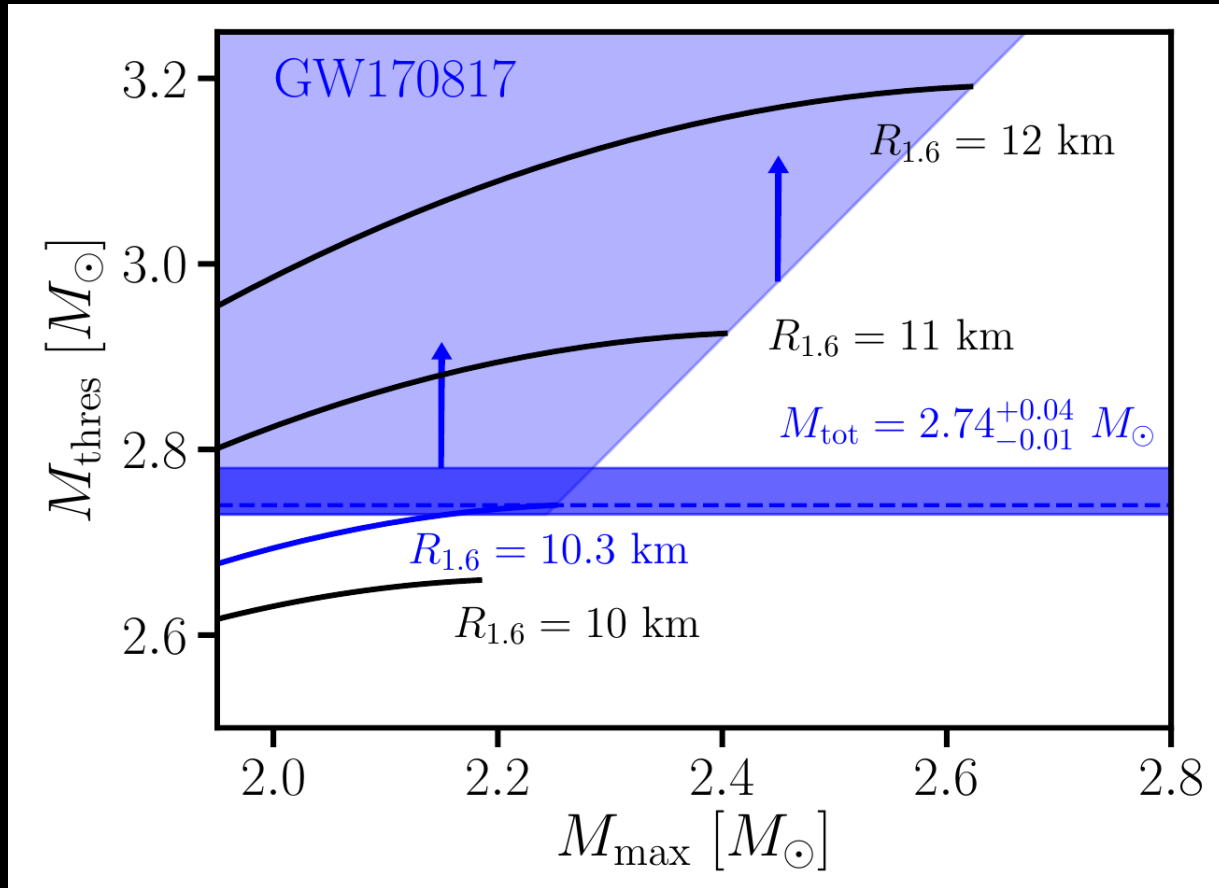
$$M_{\text{tot}}^{\text{GW170817}} \leq \left(-3.38 \frac{G M_{\text{max}}}{c^2 R_{\text{max}}} + 2.43 \right) M_{\text{max}} \leq \left(-\frac{3.38}{2.82} + 2.43 \right) \frac{1}{2.82} \frac{c^2 R_{\text{max}}}{G}$$

→ Lower limit on NS radius



$$M_{\text{thres}} = \left(-3.38 \frac{GM_{\text{max}}}{c^2 R_{\text{max}}} + 2.43 \right) M_{\text{max}}$$

$$M_{\text{thres}} \geq 1.2 M_{\text{max}}$$

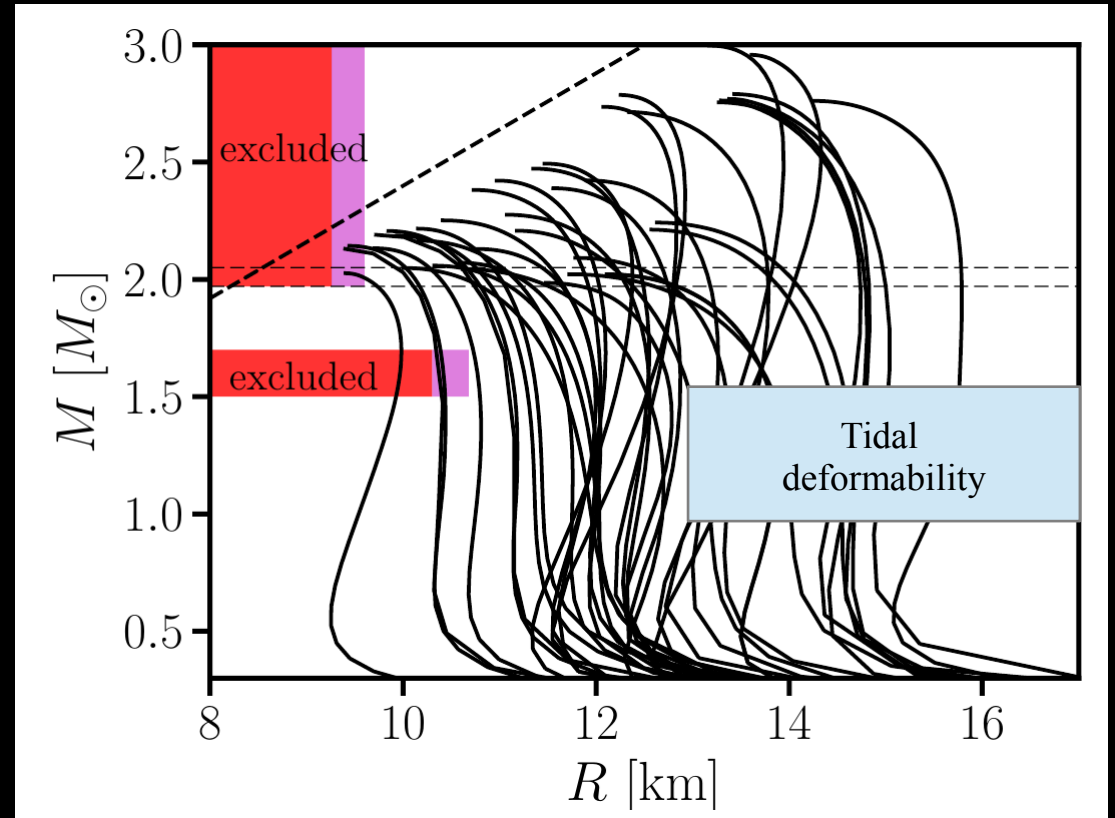


$$M_{\text{thres}} = \left(-3.6 \frac{G M_{\text{max}}}{c^2 R_{1.6}} + 2.38 \right) M_{\text{max}}$$

$$v_S = \sqrt{\frac{dP}{de}} \leq c \rightarrow M_{\text{max}} \leq \kappa R_{1.6} \Rightarrow M_{\text{thres}} \geq 1.2 M_{\text{max}}$$

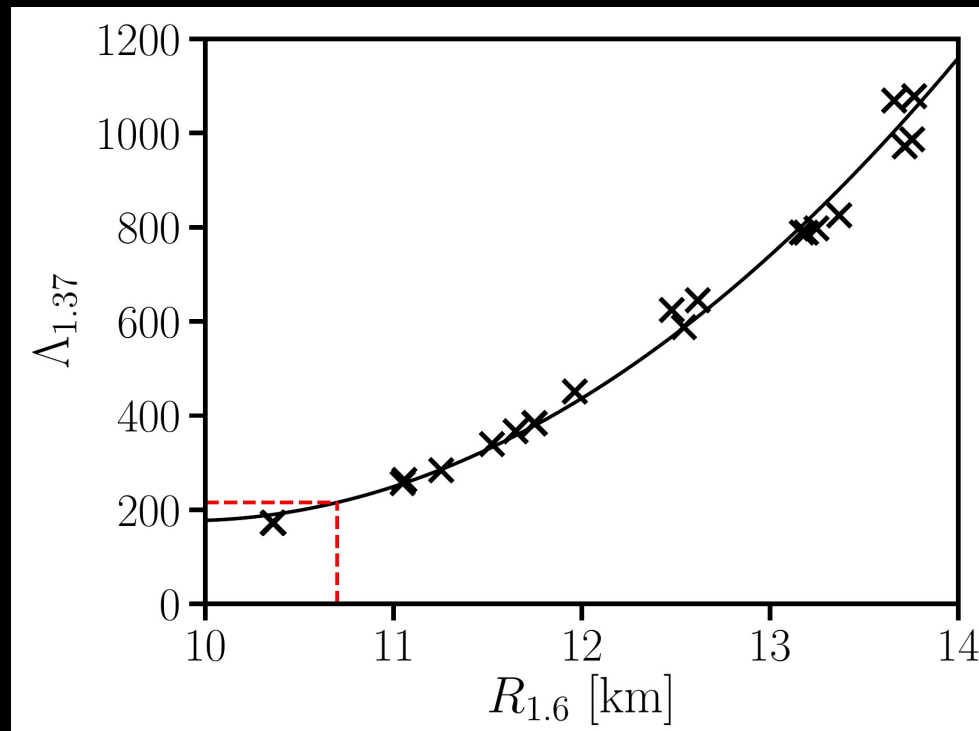
NS radius constraint from GW170817

- ▶ $R_{\text{max}} > 9.6$ km
- ▶ $R_{1.6} > 10.7$ km
- ▶ Excludes very soft nuclear matter



Bauswein et al. 2017

Radius vs. tidal deformability



Bauswein, unpubl.

- ▶ Radius and tidal deformability scale tightly → **Lambda > 210**
- ▶ Note: limit cannot be much larger because there are EoS (with somewhat larger Lambda) that do NOT result in a prompt collapse !!

(argument turned around – full EoS coverage essential !!!)

(little bit of buffer because we were very conservative in Bauswein et al 2017)

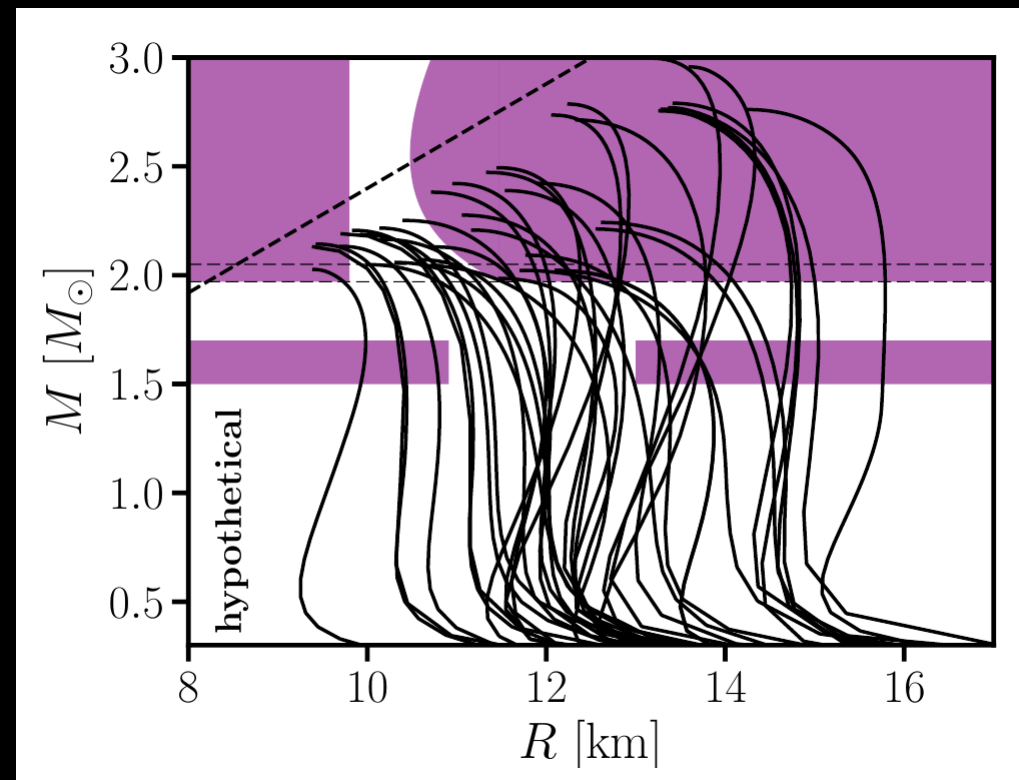
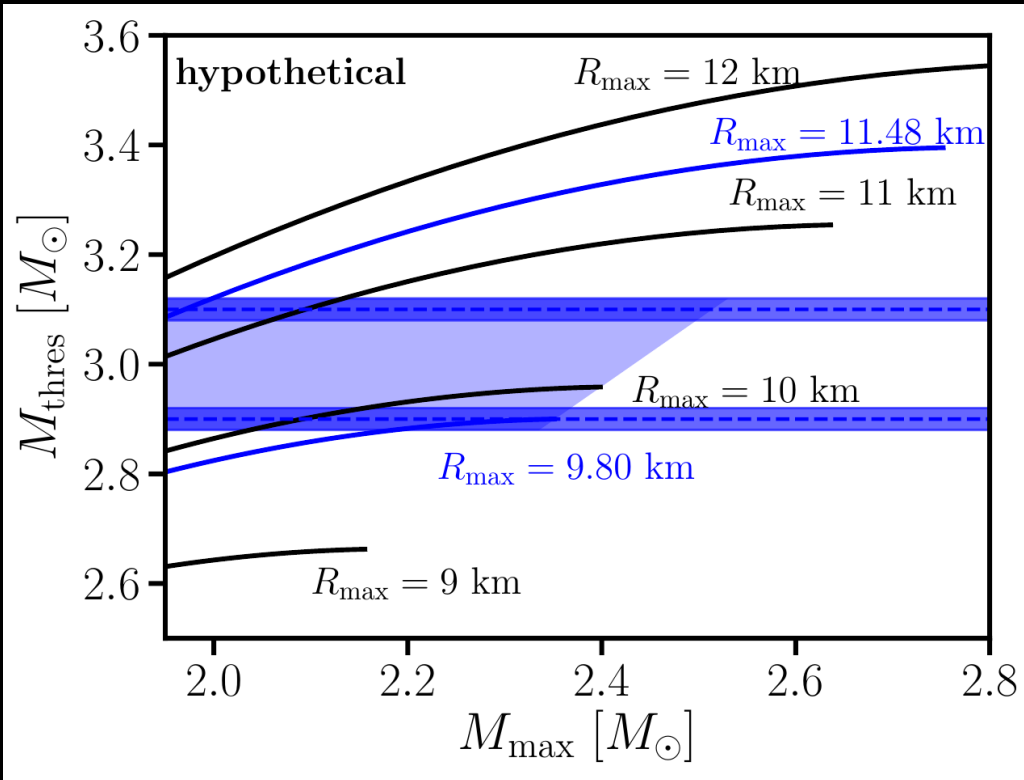
Discussion - robustness

- ▶ Binary masses well measured with high confidence error bar
- ▶ Clearly defined working hypothesis: delayed collapse
 - testable by refined emission models
 - as more events are observed more robust distinction
- ▶ Very conservative estimate, errors can be quantified
- ▶ Empirical relation can be tested by more elaborated simulations (but unlikely that MHD or neutrinos can have strong impact on M_{thres})
- ▶ Confirmed by semi-analytic collapse model
- ▶ Low-SNR constraint !!!

Future

- ▶ Any new detection can be employed if it allows distinction between prompt/delayed collapse
- ▶ With more events in the future our comprehension of em counterparts will grow → more robust discrimination of prompt/delayed collapse events
- ▶ Low-SNR detections sufficient !!! → that's the potential for the future
 - we don't need louder events, but more
 - complimentary to existing ideas for EoS constraints

Future detections (hypothetical discussion)



- as more events are observed, bands converge to true M_{thres}
- prompt collapse constrains M_{\max} from above

Future: Maximum mass

- ▶ Empirical relation

$$M_{\text{thres}} = \left(-3.6 \frac{G M_{\text{max}}}{c^2 R_{1.6}} + 2.38 \right) M_{\text{max}}$$

- ▶ Sooner or later we'll know $R_{1.6}$ (e.g. from postmerger) and M_{thres} (from several events – through presence/absence of postmerger GW emission or em counterpart)

=> direct inversion to get precise estimate of M_{max}

(see also current estimates e.g. by Margalit & Metzger, Rezzolla et al, Ruiz & Shapiro, Shibata et al., ...)

Postmerger GW emission*

(dominant frequency of postmerger phase)

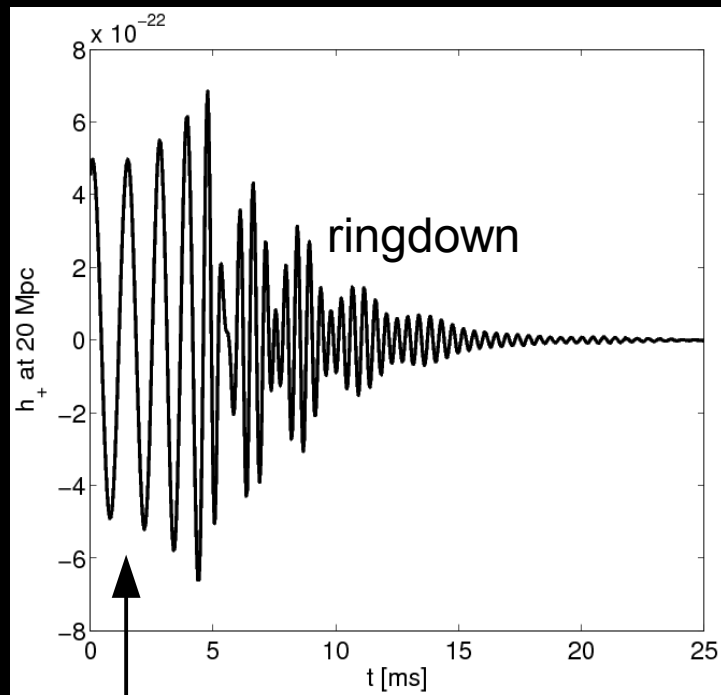
→ determine properties of EoS/NSs

→ complementary to inspiral

* not detected for GW170817 – expected for current sensitivity and $d=40$ Mpc
(Abbott et al. 2017)

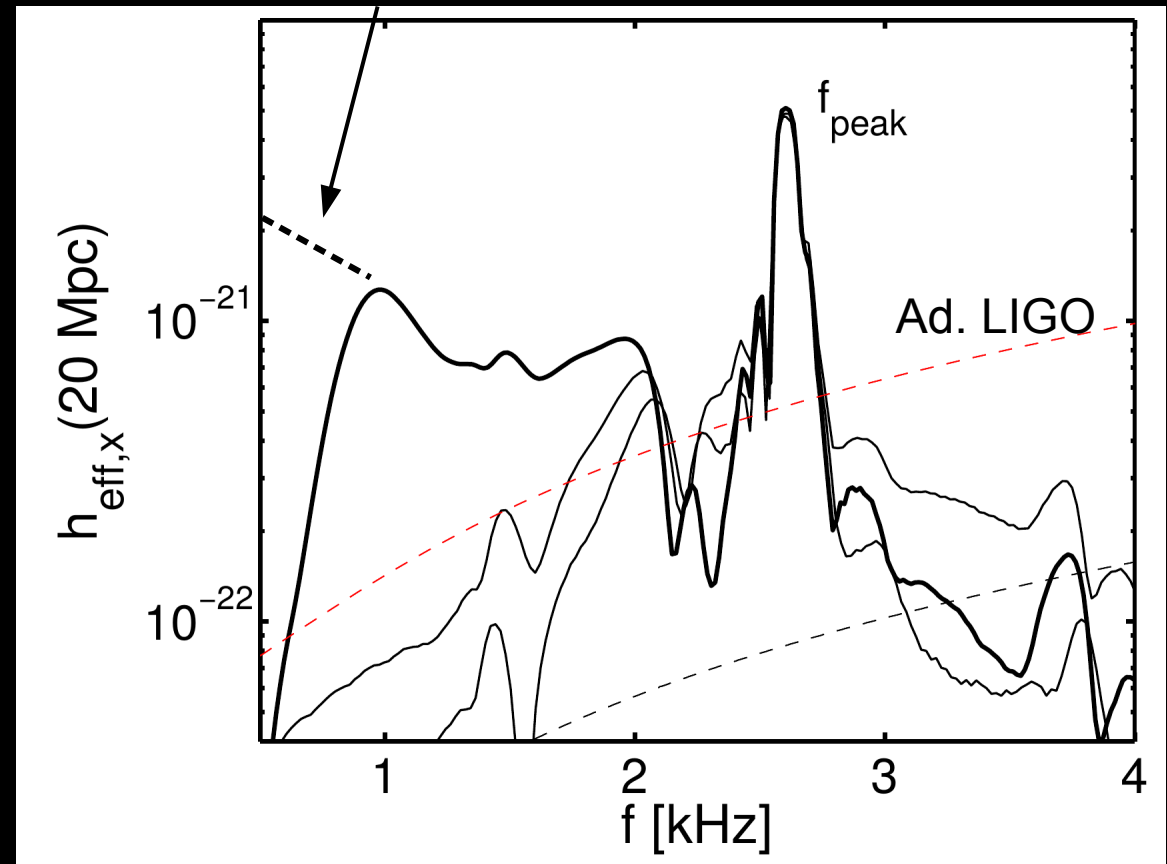
Postmerger

1.35-1.35 M_{sun} , 20 Mpc



Earlier inspiral

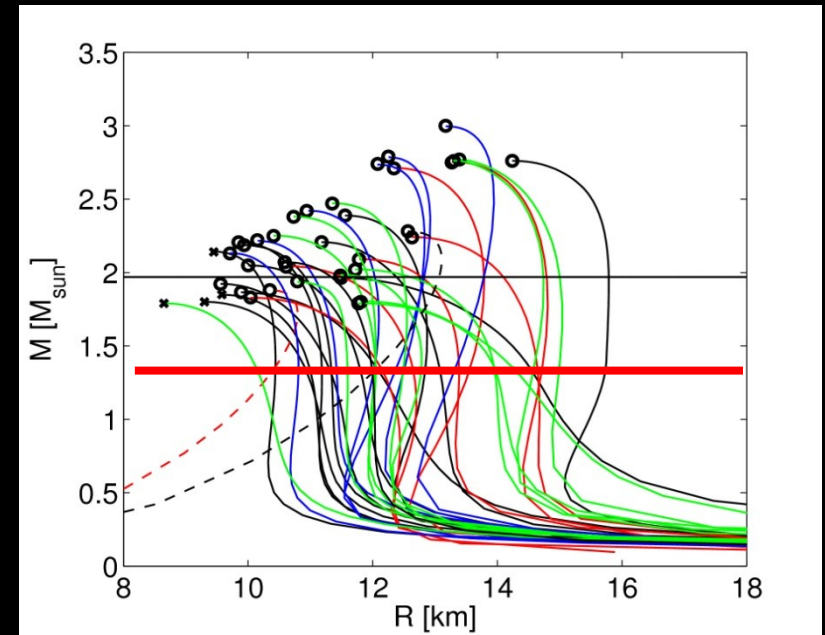
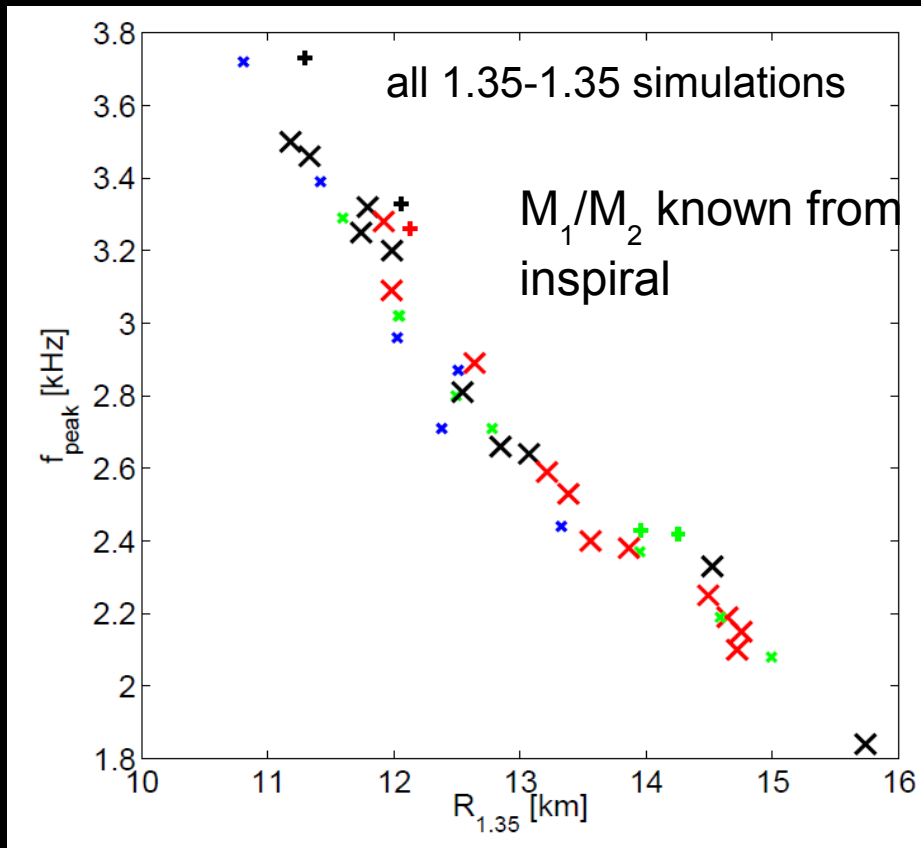
EoS



Dominant postmerger oscillation frequency f_{peak}

Very characteristic (robust feature in all models)

Gravitational waves – EoS survey



characterize EoS by radius of nonrotating NS with $1.35 M_{\text{sun}}$

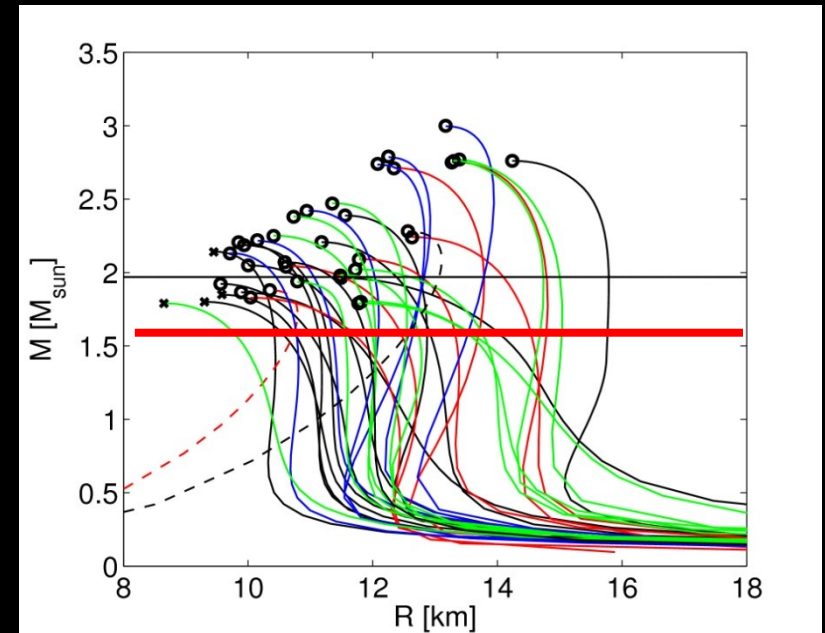
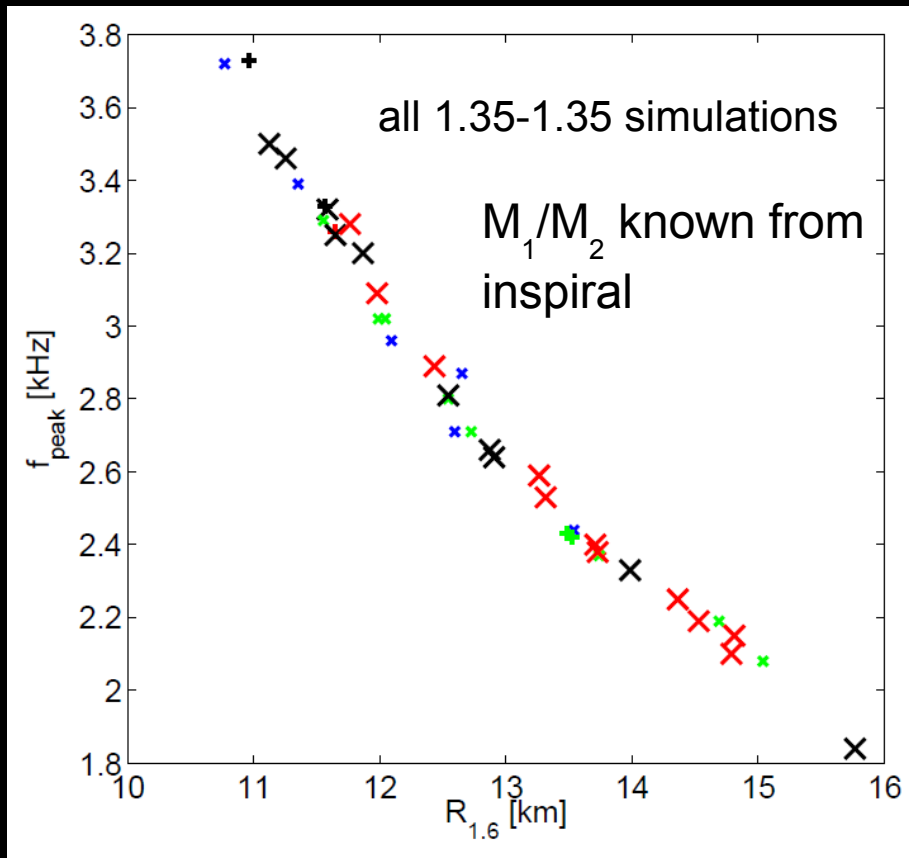
Bauswein et al. 2012

Pure TOV/EoS property => **Radius measurement** via f_{peak}

Here only 1.35-1.35 M_{sun} mergers (binary masses measurable) – similar relations exist for other fixed binary setups !!!

~ 40 different NS EoSs

Gravitational waves – EoS survey



characterize EoS by radius of nonrotating NS with $1.6 M_{\text{sun}}$

Bauswein et al. 2012

Pure TOV/EoS property \Rightarrow **Radius measurement** via f_{peak}

Smaller scatter in empirical relation (< 200 m) \rightarrow smaller error in radius measurement

Note: R of $1.6 M_{\text{sun}}$ NS scales with f_{peak} from 1.35 - $1.35 M_{\text{sun}}$ mergers (density regimes comparable)

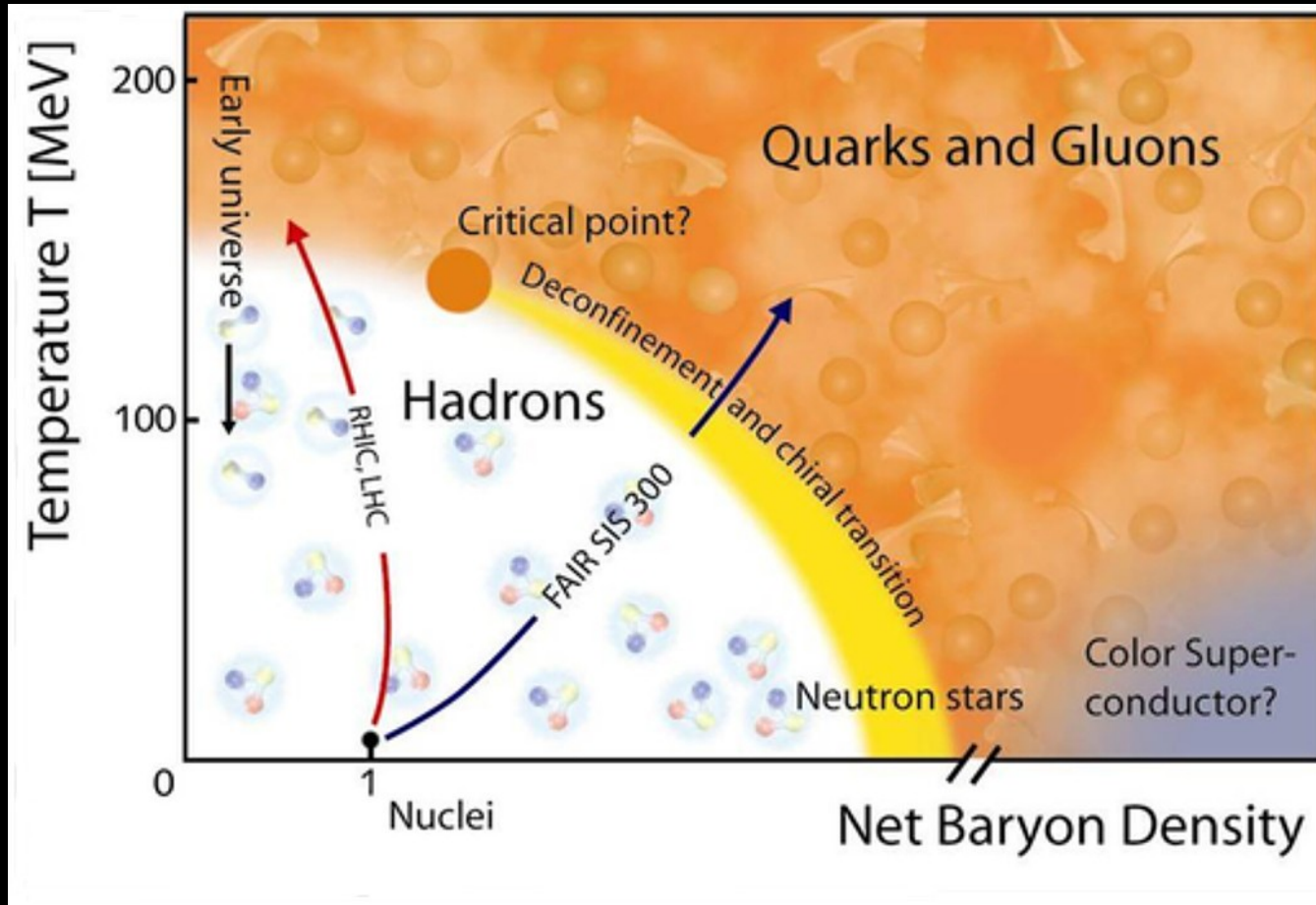
GW data analysis: Clark et al 2014, Clark et al 2016, Chatziioannou et al 2017, ...

\rightarrow detectable at a few 10 Mpc

Observable signature of (QCD) phase transition

Phase diagram of matter

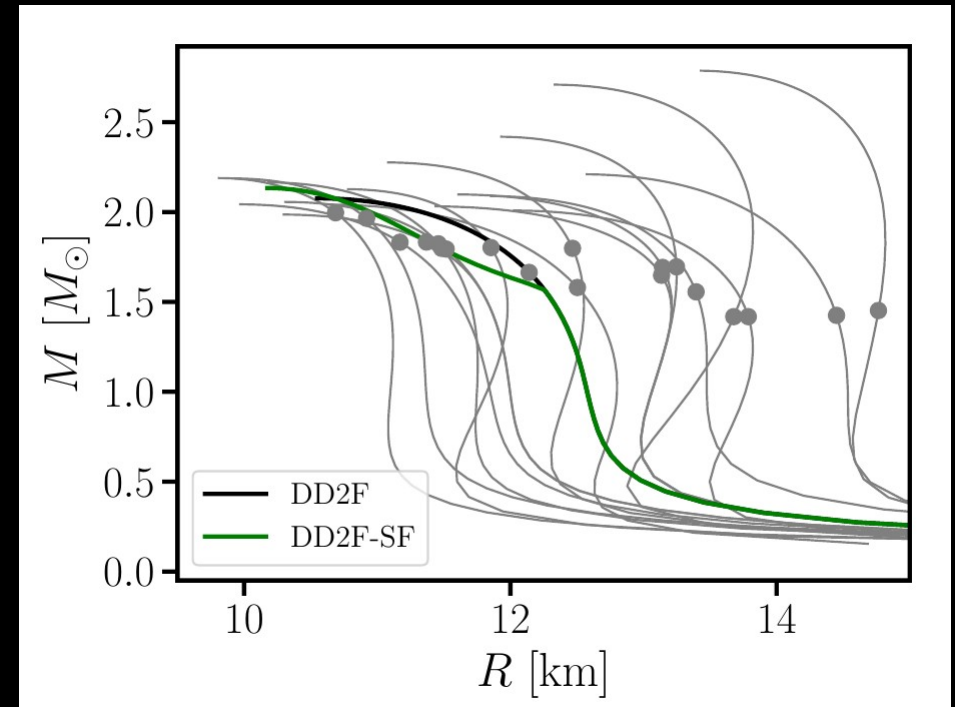
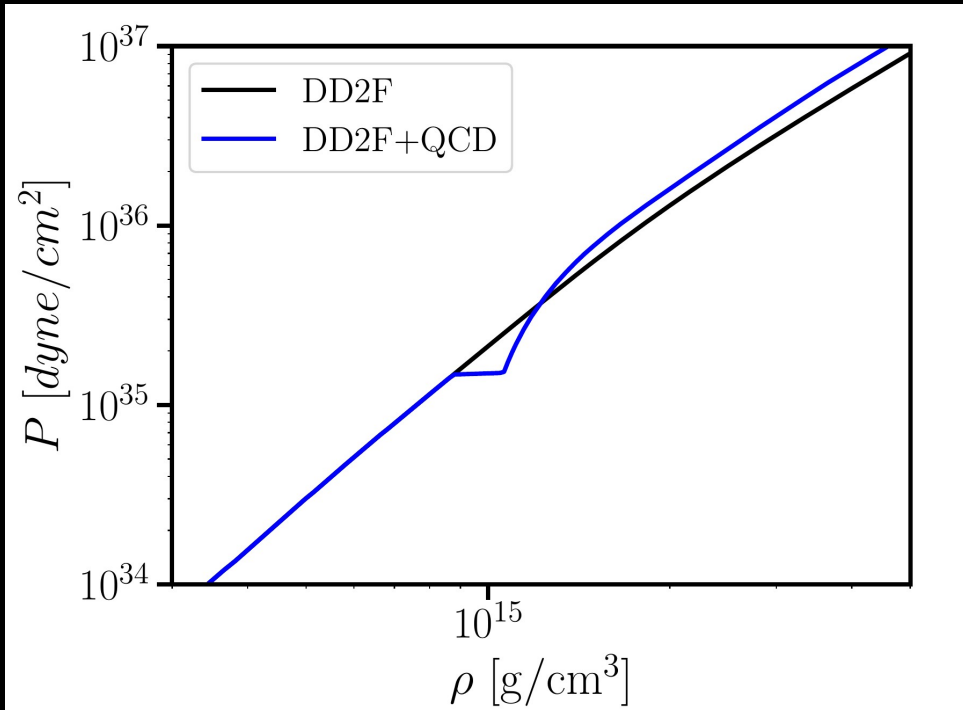
GSI/FAIR



Does the phase transition to quark-gluon plasma occur (already) in neutron stars or only at higher densities ?

EoS with 1st-order phase transition to quark matter

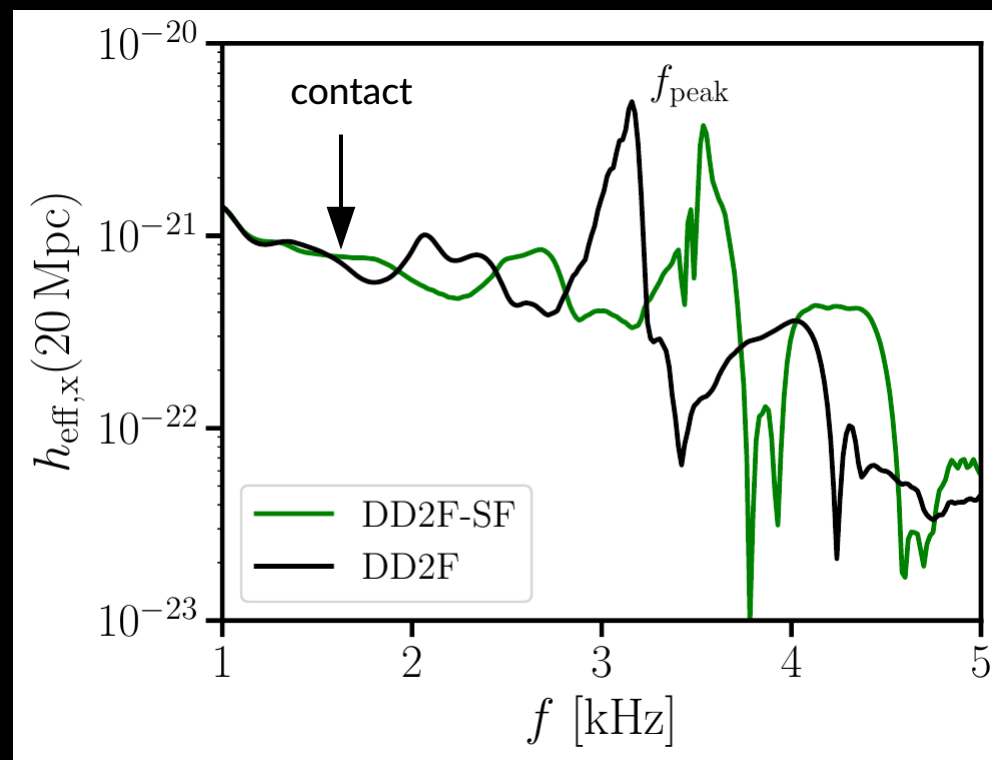
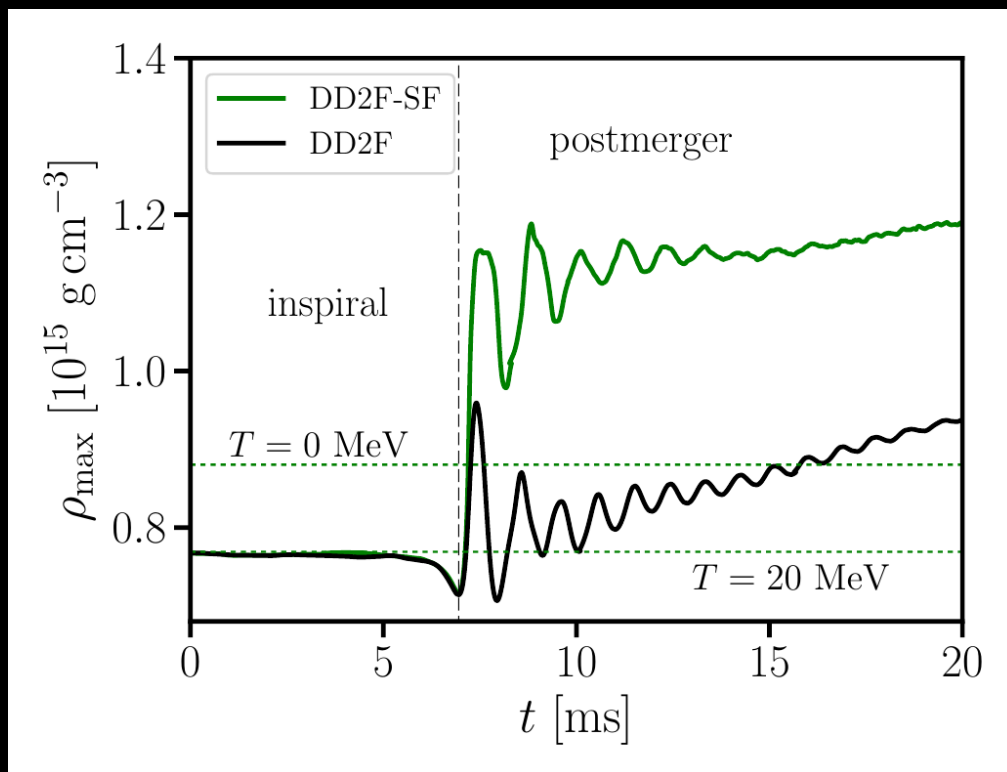
Bauswein et al. 2018



- ▶ EoS from Fischer et al. 2018 – as one example for an EoS with a strong 1st-order phase transition to deconfined quarks
- ▶ Difficult to measure transition in mergers through inspiral: Λ very small, high mass star probably less frequent

Merger simulations

► GW spectrum 1.35-1.35 Msun



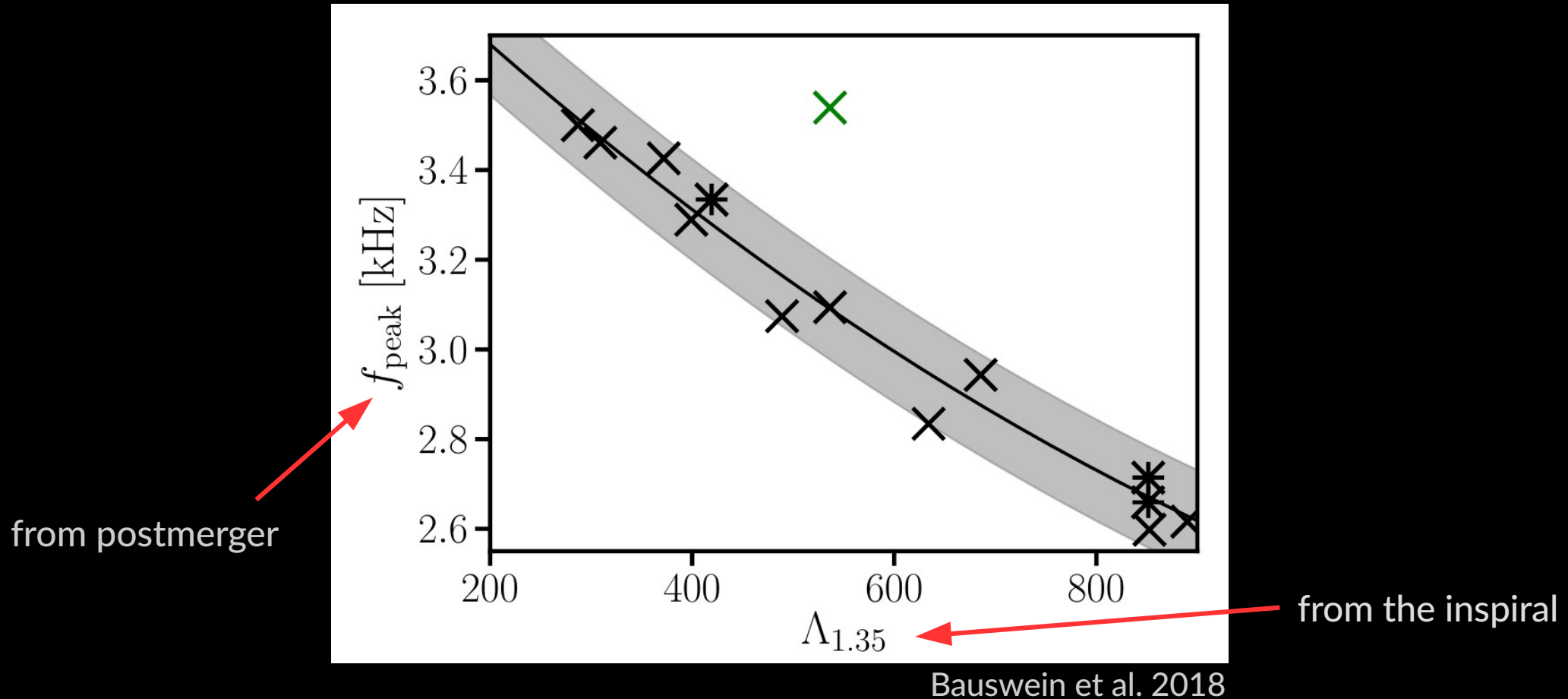
Bauswein et al. 2018

But: a high frequency on its own may not yet be characteristic for a phase transition

→ unambiguous signature

(→ show that all purely baryonic EoS behave differently)

Signature of 1st order phase transition

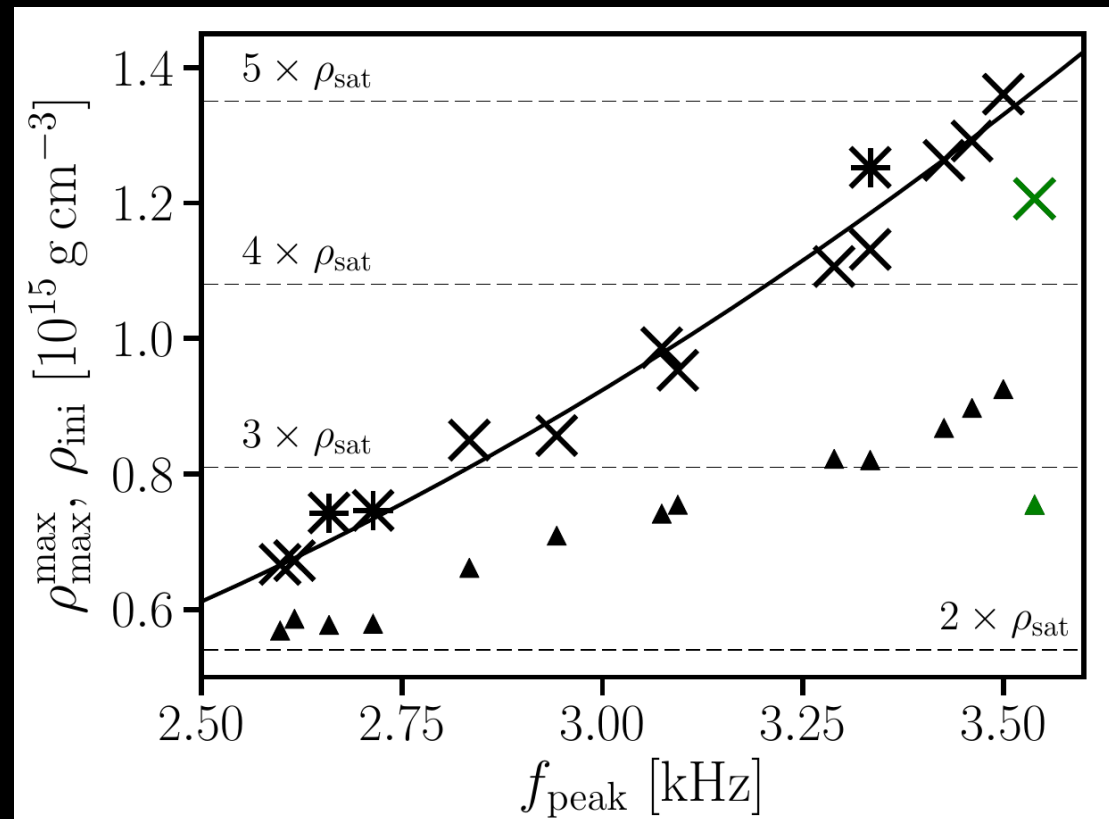


- ▶ Tidal deformability measurable from inspiral to within 100-200 (Adv. Ligo design)
- ▶ Postmerger frequency measurable to within a few 10 Hz @ a few 10 Mpc (either Adv. Ligo or upgrade)
- ▶ Important: “all” purely hadronic EoSs (including hyperonic EoS) follow f_{peak} -Lambda relation → deviation characteristic for strong 1st order phase transition

Discussion

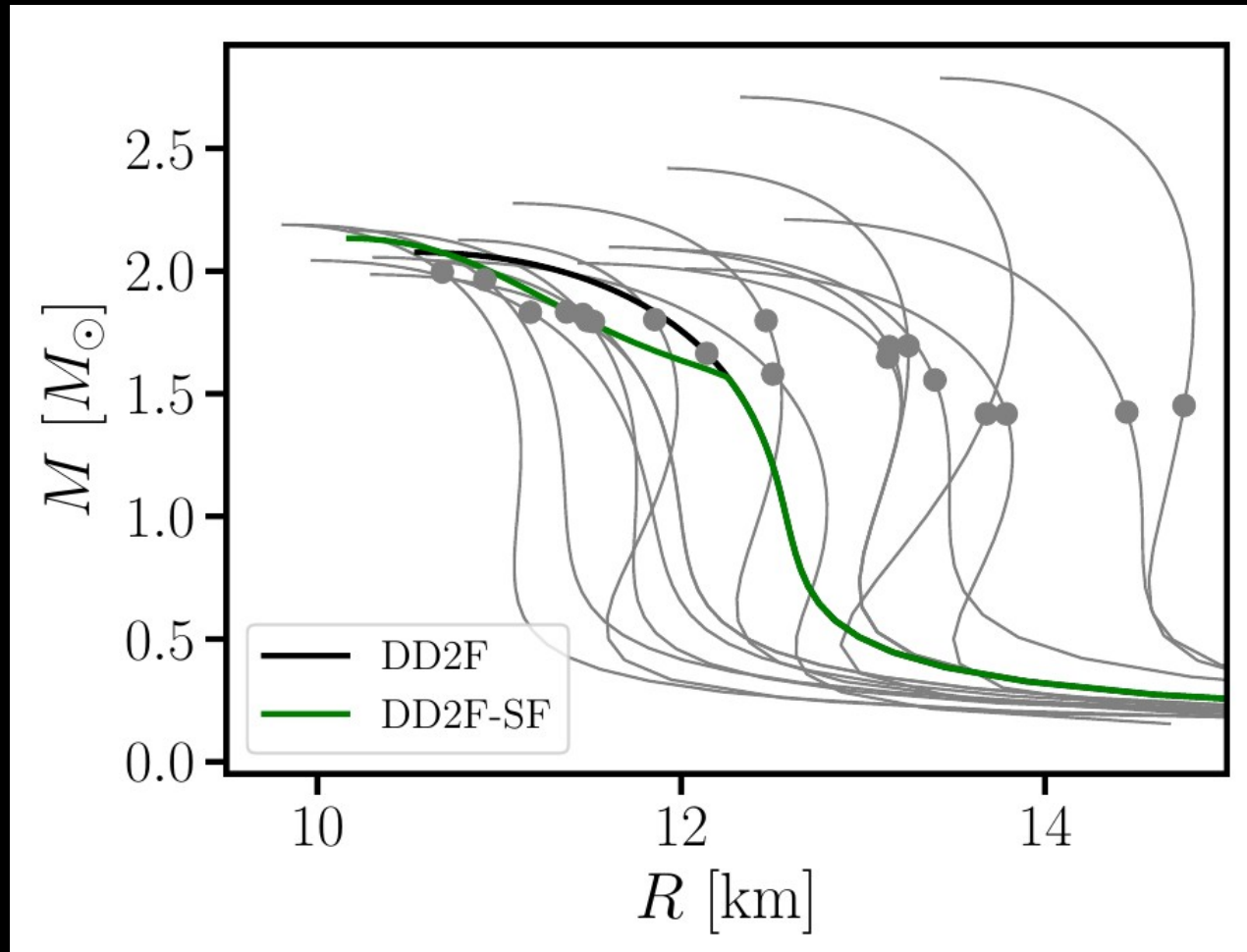
- ▶ Consistency with f_{peak} -Lambda relation points to
 - purely baryonic EoS
 - (or an at most weak phase transition \rightarrow no strong compactification)in the tested (!) density regime
- ▶ f_{peak} also determines maximum density in postmerger remnant
- ▶ postmerger GW emission provides complimentary information to inspiral \rightarrow probes higher density regime

Bauswein et al. 2018



Probed densities / NS masses

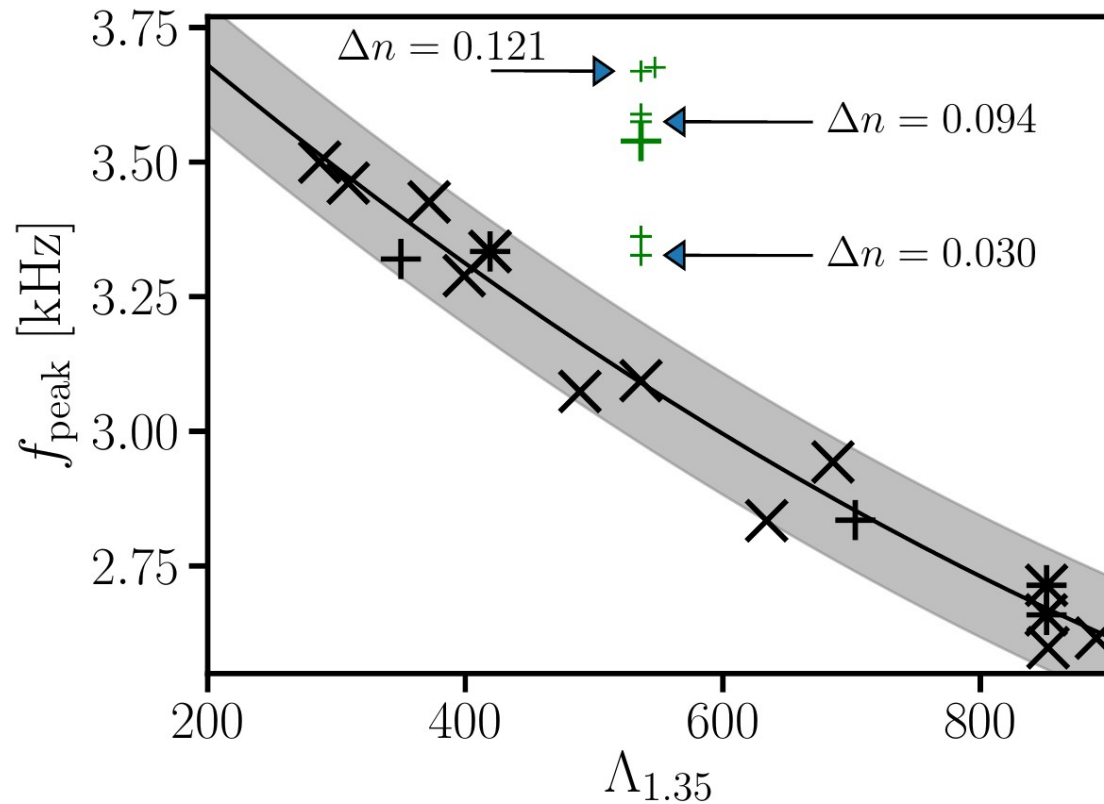
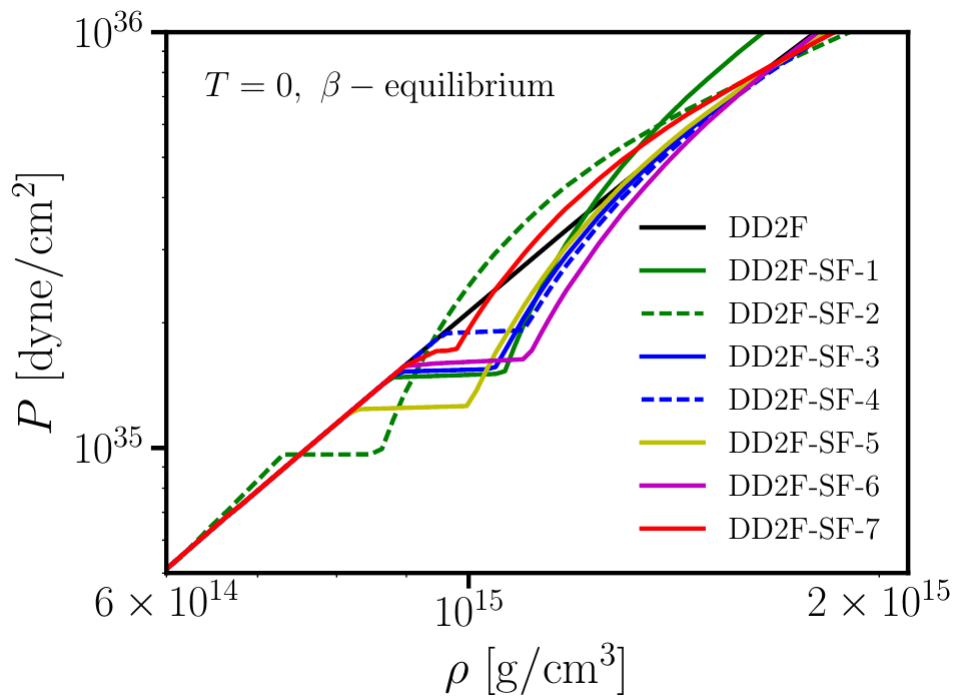
- ▶ Dots: NS mass with central density = maximum density during early postmerger evolution



For 1.35-1.35 M_{sun} merger – higher binary masses probe higher densities / NS masses

More models

- ▶ Larger density jump \rightarrow stronger compactification \rightarrow more significant increase of f_{peak} (keeping other EoS parameters fixed)



Conclusions

- ▶ NS radius must be larger than 10.7 km (very robust and conservative)
- ▶ More stringent constraints from future detections
- ▶ NS radius measurable from dominant postmerger frequency
- ▶ Explicitly shown by GW data analysis
- ▶ Threshold binary mass for prompt collapse → maximum mass M_{\max}
- ▶ Strong 1st order phase transitions leave characteristic imprint on GW (postmerger frequency higher than expected from inspiral)
- ▶ Complementarity of inspiral and postmerger phase → postmerger probes higher density regime